

Fragmentation of SQM in astrophysical events



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Strangelets in CR: some hints and the AMS02 experiment

Centauro events (Bjorken & McLerran 1979)

ET event (Ichimura et al. 1993)

...

Last report

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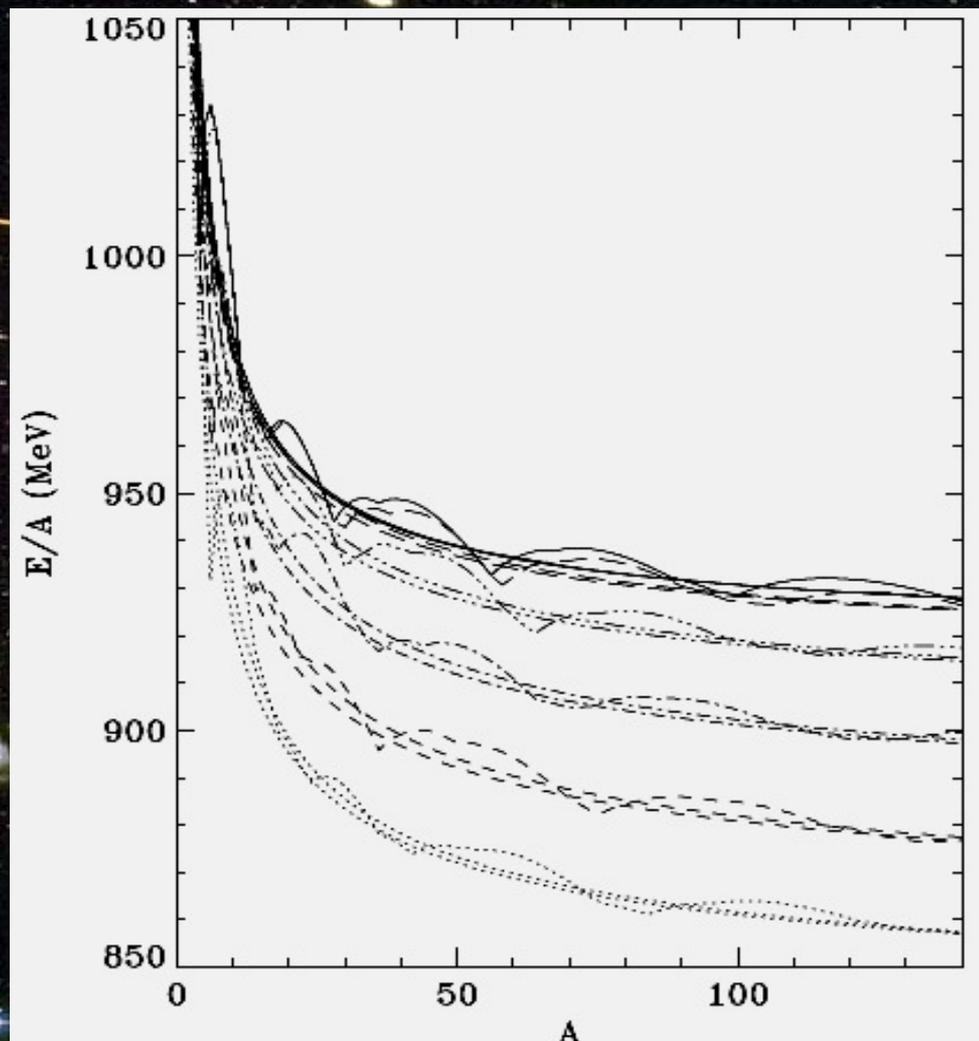
Observation of a rare cosmic ray event at mountain altitude

Basudhara Basu^{a, b}, Sibaji Raha^{a, b, ,}, Swapan K. Saha^{a, b}, Sukumar Biswas^{c, 1},
Sandhya Dey^b, Atanu Maulik^b, Amal Mazumdar^d, Satyajit Saha^e, Debapriyo Syam^f

The problem: what happens when SQM is ejected?

- SS mergers
- Type II supernovae

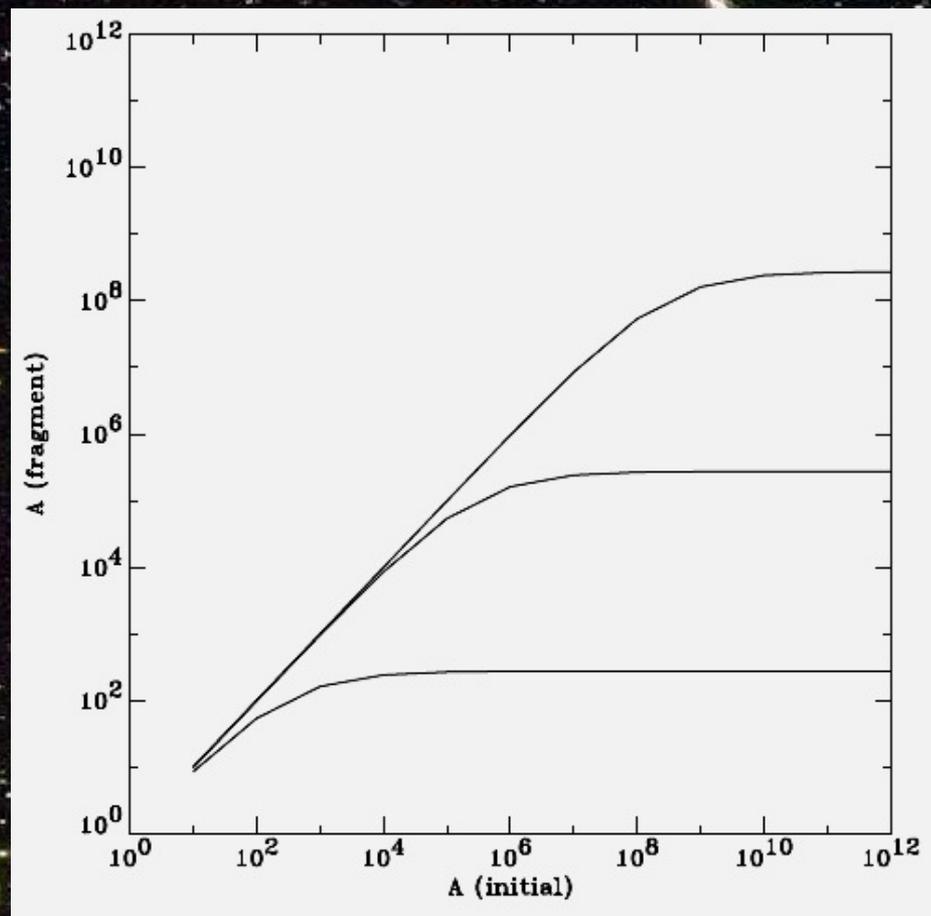
Tentative answer:
(almost) nothing



Madsen (2002):

A number for which gravity equals surface energy $\sim 10^{38}$

Further collisions between chunks in orbit fragment into strangelets down to $\sim 10^3$



Horvath & Vucetich (1998): fragmentation of strangelets on their way out a massive (SNII) oxygen envelope

Spallation cross-section of a fragment Δ with oxygen (Boyd & Saito 1993)

$$\sigma(\Delta) = \sigma_0 \left(\frac{m}{m_0} \right)^{2/3} \exp(-\Delta/\Delta_0)$$

Heavy strangelets colliding with oxygen will lose mass and baryon number



$$\frac{dm}{dt} = - \int d\Delta n(t) \sigma(\Delta) v$$

$$m \frac{dv}{dt} = -C \frac{\pi}{2} m_{\text{ox}} n(t) v^2 R_0^2 \left(\frac{m}{m_0} \right)^{2/3} + \dot{m} \xi v$$

$$n(t) = n_0 \exp(-t/\tau_{\text{exp}})$$

Fowler & Hoyle
(1964)

Whenever $D = (\pi C m_{ox} R_0^2) / (2\Delta_0 \sigma_0) + \xi > 1$, strangelets loose mass and kinetic energy, freezing out at

$$E_{i,ox} \left(\frac{m}{m_i} \right)^{2(D-1)} = \frac{\Delta_0}{m_p} E_b.$$

$$\frac{m_F}{m_i} = \left(\frac{\Delta_0}{m_p} \frac{E_b}{E_{i,ox}} \right)^{1/2(D-1)}$$

Scaling law for the escaping strangelet mass

m_F

Strangelet total mass must be small, since SNII are frequent. Herzog & Ropke (2011) showed that instabilities leading to turbulent fast deflagrations (Horvath & Benvenuto (1988); Horvath (2010)) do their job, but it is unclear whether any SQM is ejected at all...

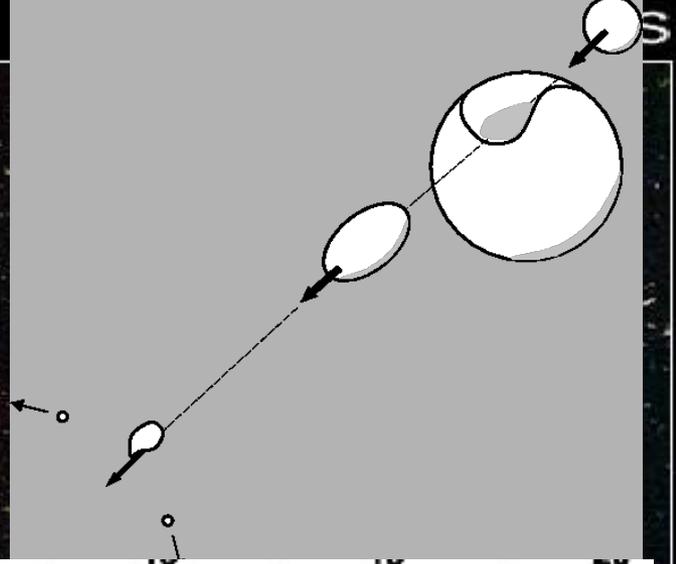
Abrasion-ablation model of spallation (GEANT4 and others)

$$\Delta_{abr} = F A_p [1 - \exp(-C_T/\lambda)]$$

Fraction of projectile in interaction

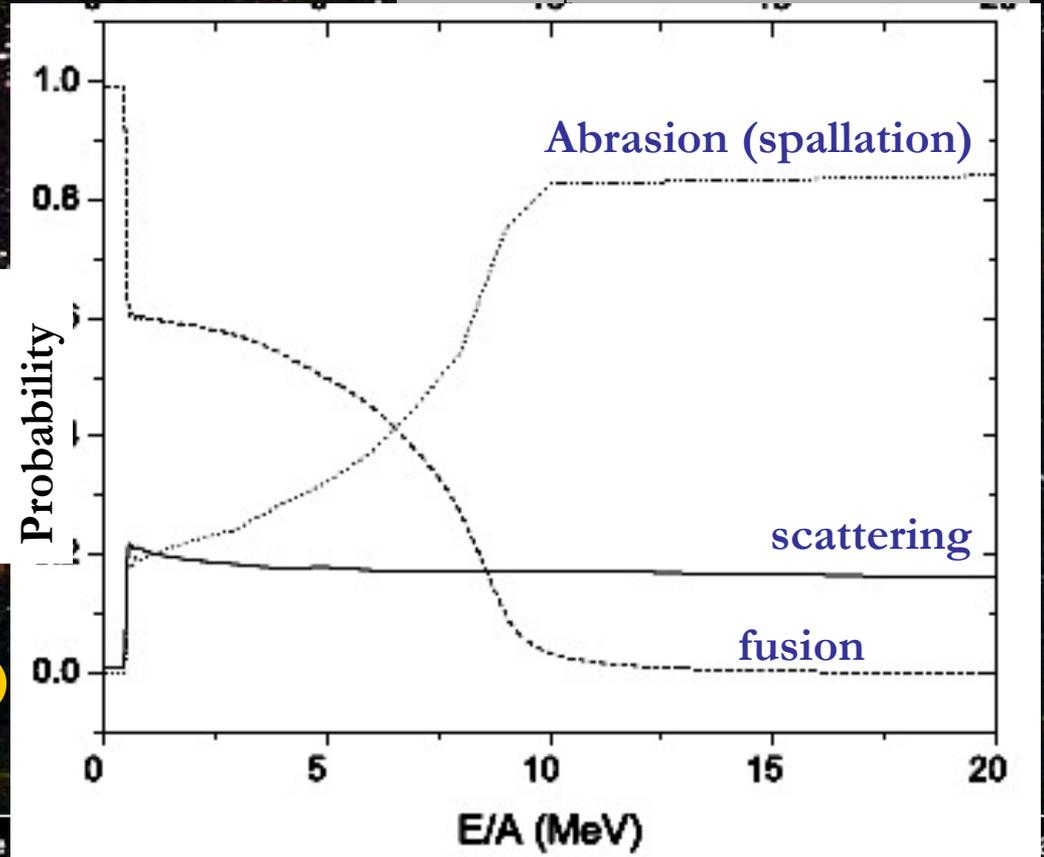
Baryon number of projectile

Mean free path of N-N interaction



Strangelet + nitrogen
w/GEANT4: few
tens of A stripped

(Paulucci & Horvath JPG 2011)



On the other hand, in mergers free expansion + cooling of ejected matter should be relevant to fragmentation, independently of collisions

Statistical Multifragmentation Model (SMM)

Needs T (excitation) $> \sim$ binding

Partition function of a fragment with baryon number A

$$\omega_A = V \left(\frac{mTA}{2\pi} \right)^{3/2} e^{-f_A/T}$$

Internal free energy of the fragment A

$$f_A = -WA + \sigma A^{2/3} + CA^{1/3}$$

Binding energy of
SQM in bulk

Surface

Curvature

In the grand canonical ensemble
the pressure is

Fourier transform of the
partition function

$$p(T, \mu) = T \lim_{V \rightarrow \infty} \frac{\ln Z(V, T, \mu)}{V}$$

The pressures of each phase are obtained from the singularities of
the isobaric partition function

$$p_g(T, \mu) = T \left(\frac{mT}{2\pi} \right)^{3/2} \left\{ z_1 e^{\frac{\mu - bp_g}{T}} + \sum_{A=2}^{\infty} A^{3/2} e^{[(v - bp_g)A - \sigma A^{2/3} - CA^{1/3}]/T} \right\}$$

$$p_l(T, \mu) = \frac{v}{b}$$

with

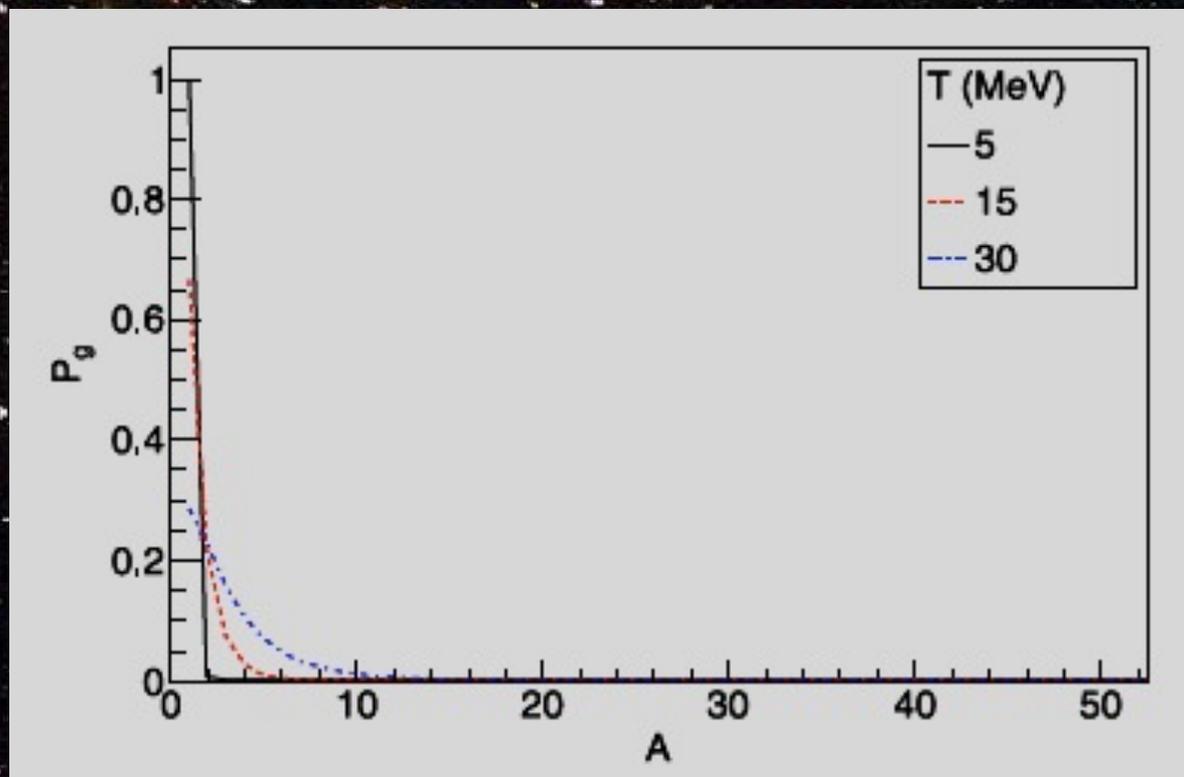
$$v = \mu + W$$

The fragmentation spectrum (assuming chemical equilibrium at the fragmentation point) is

$$\mathcal{P}_g(A) = \frac{\partial}{\partial \mu_A} p_g = \left(\frac{m_0 T}{2\pi} \right)^{3/2} A^{3/2} e^{[(\mu+W - bp_g)A - \sigma A^{2/3} - CA^{1/3}]/T}$$

Repulsive forces included
in a VdW approximation

The crisis: $A_{\text{peak}} \ll 1$



$B^{1/4} = 145$ MeV, $m_s = 100$ MeV, and $\Delta = 50$ MeV.

The role of the bag: melting the vacuum

A fraction of the vacuum energy is being used to create the surfaces, hence there is an energy density difference between both phases

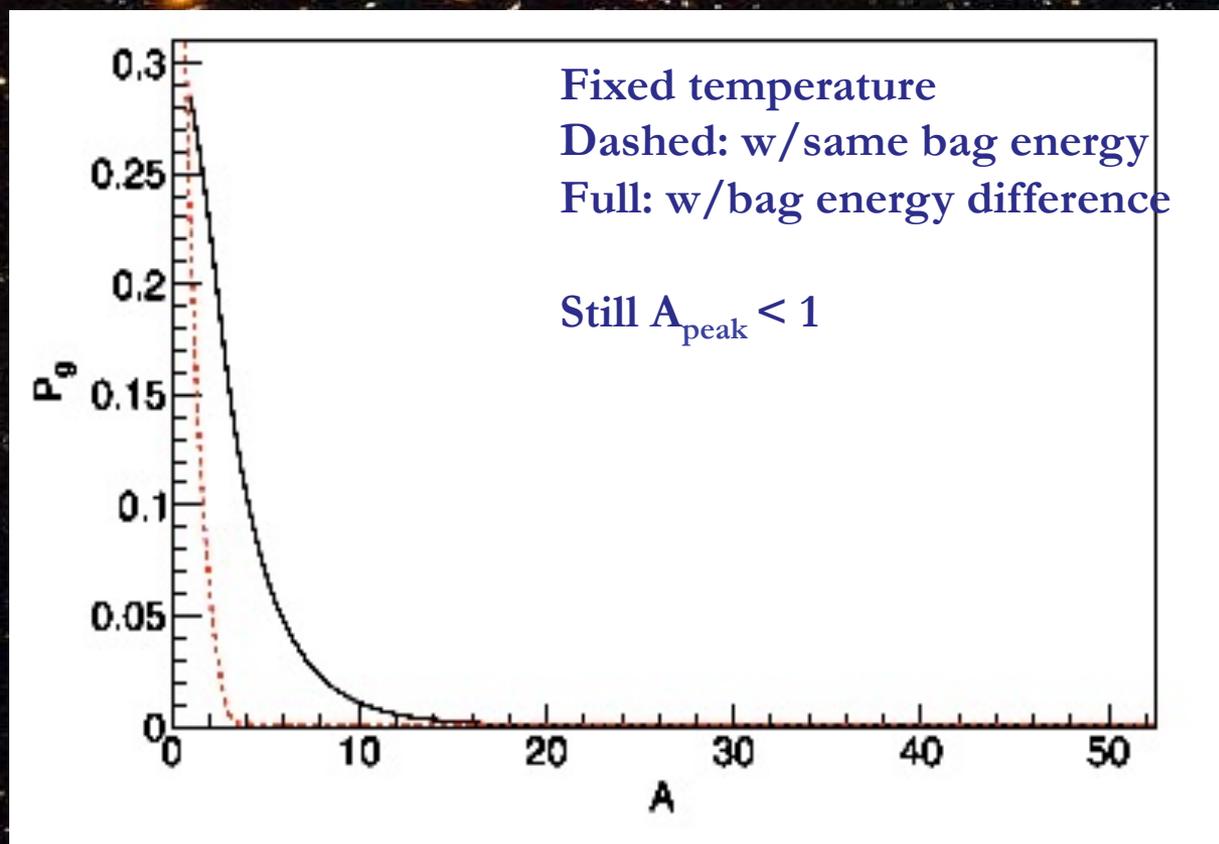
$$W_l = W_0 + Bv_{liq},$$

$$W_g = W_0 + Bv_{gas}.$$

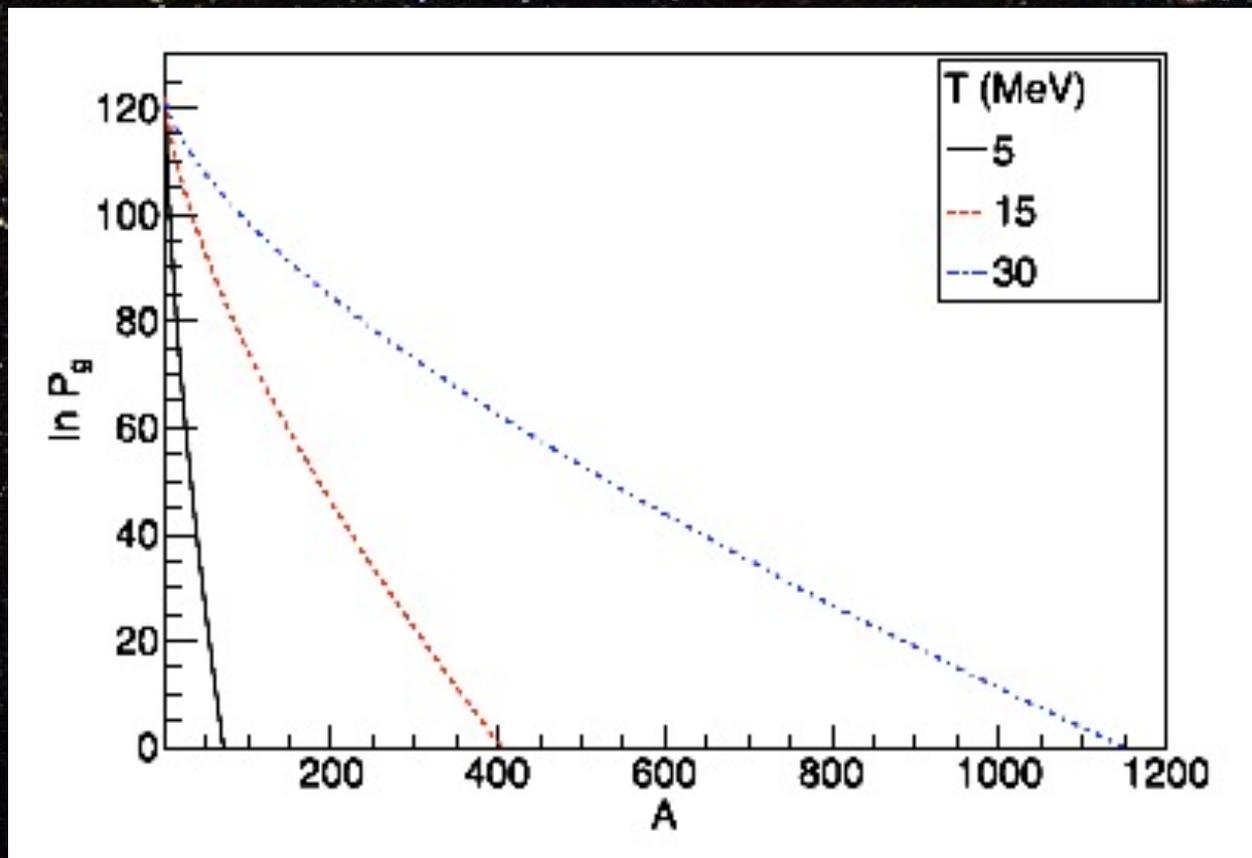
In the coexistence phase

$$[\nu - bp_g^*]A = [\nu_g - \nu_l^*]A = [\mu_g + W_g - \mu_l - W_l]A = [B(v_{gas} - v_{liq})]A.$$

The crisis does not go away:

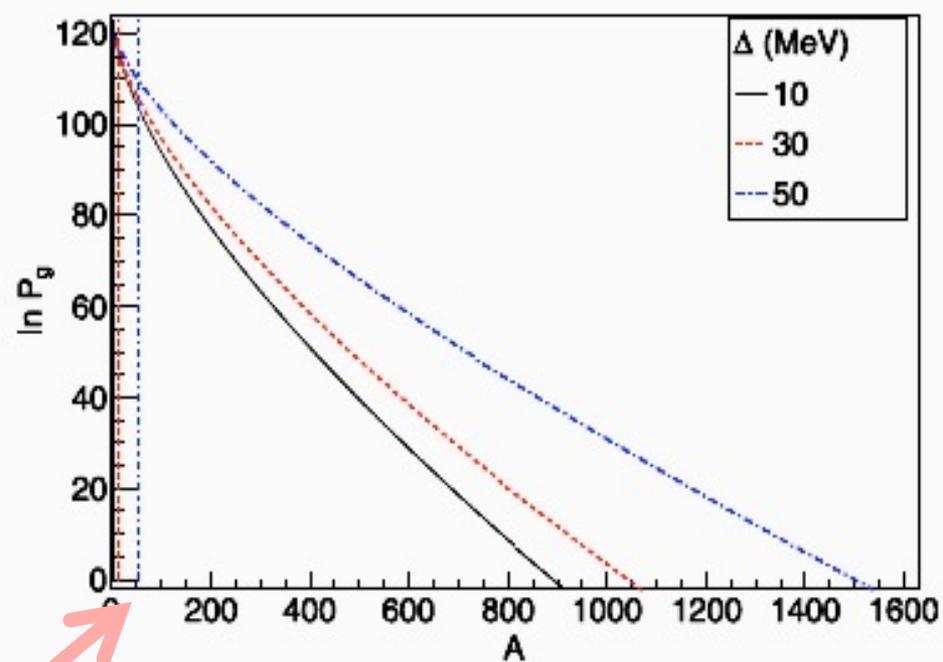
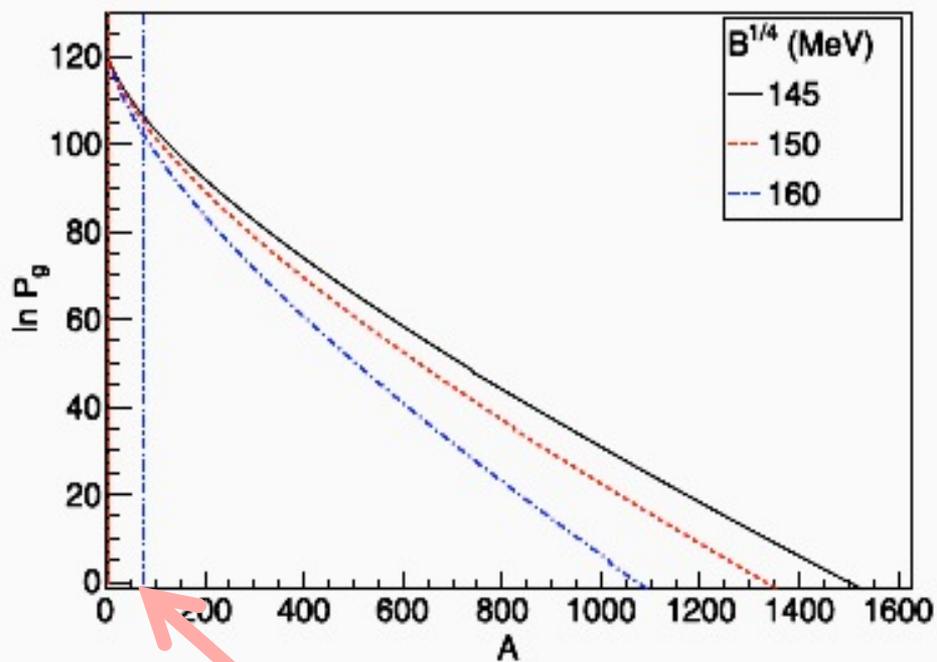


$$B^{1/4} = 145 \text{ MeV}, m_s = 100 \text{ MeV}, \text{ and } \Delta = 50 \text{ MeV}$$



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Probability distribution for a fixed $10^{-4}M_{\odot}$ ejected mass of SQM

Effects of B and pairing Δ 

Note the self-consistent determination of A_{\min}

Conclusions

- SQM does not like to fragment, since most of the fragmentation spectrum peaks at $A_{\text{peak}} \ll 1$
(Biswas *et al.* (PLB 2012) just present the normalized probabilities but do not mention this problem, in their work $m_s = 0$)
- Strangelets with $A > 1$ in the exponential part of the distribution subject to decay into ordinary hadrons because they are smaller than A_{min}

- A very small fraction of the total ejected SQM would produce strangelets → the flux should be minuscule in anomalous CR detection experiments

$$F \approx 5 \times 10^5 (\text{m}^2 \text{ y sterad})^{-1} \times R_{-4} \times M_{-2} \times V_{100}^{-1} \times t_7,$$

Rate of mergers

There is an very small “efficiency” factor $\sim 10^{-5}$ to be inserted here

→ Bauswein et al. (2010) : the ejection may even be zero (independently of efficiency)

- It is possible that fragmentation happens out of (chemical and mechanical?) equilibrium

A vast field of galaxies, including spiral, elliptical, and irregular shapes, scattered across a dark cosmic background. The galaxies vary in size, color, and orientation, representing a wide range of stellar populations and evolutionary stages.

Danken !