

# INFERRING NEUTRON STARS CRUST PROPERTIES

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with A. Turlione (PhD student) and J. Pons  
arXiv: 1309.3909

- Motivation
- Five sources & open issues
- Theoretical approach
- Envelope models
- "Crustal coolers"
- Beyond crustal cooling --> (?)
- Conclusions



Saturday, September 27, 2014



# LMXB

Quiescent state  $\rightarrow$  little accretion  $L_x \sim 10^{\{34\}}$  erg/s

Active phase  $\rightarrow$  accretion rises  $L_x \sim 10^{\{36-39\}}$  erg/s



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of the envelope

thermonuclear fusion  
Type I X-Ray Bursts



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pycnonuclear reactions  
e-captures, n-emission

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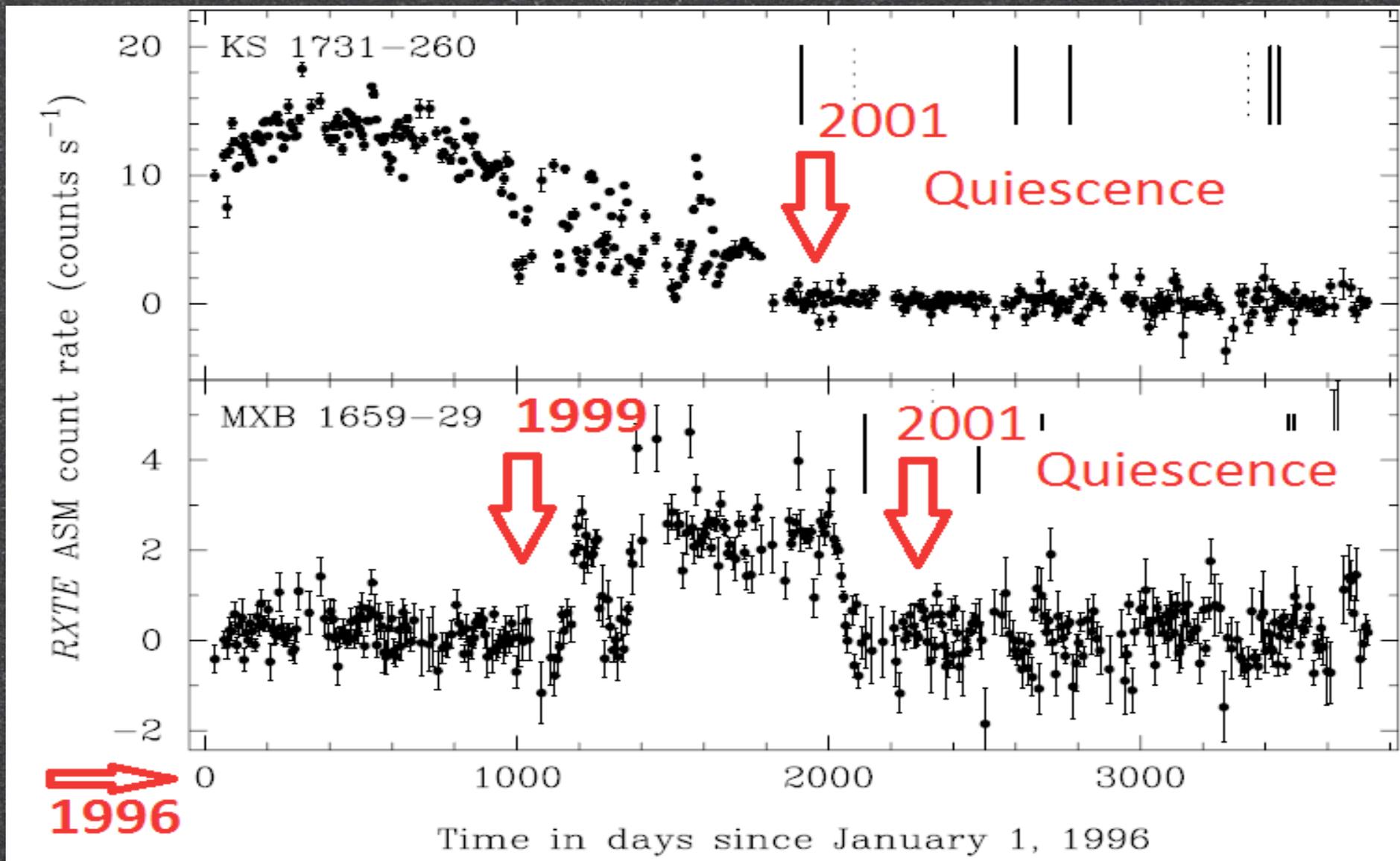
Quasi-persistent sources (accretion  $\sim$  1-10 yrs)

Accretion stops  $\rightarrow$  X-ray quiescent emission

$\rightarrow$  **CRUSTAL COOLING**

Brown et al (98, 09), Colpi et al (01), Rutledge (02), Shternin et al (07)

# NSs going into quiescence



Cackett et al (2006)

# Five sources detected

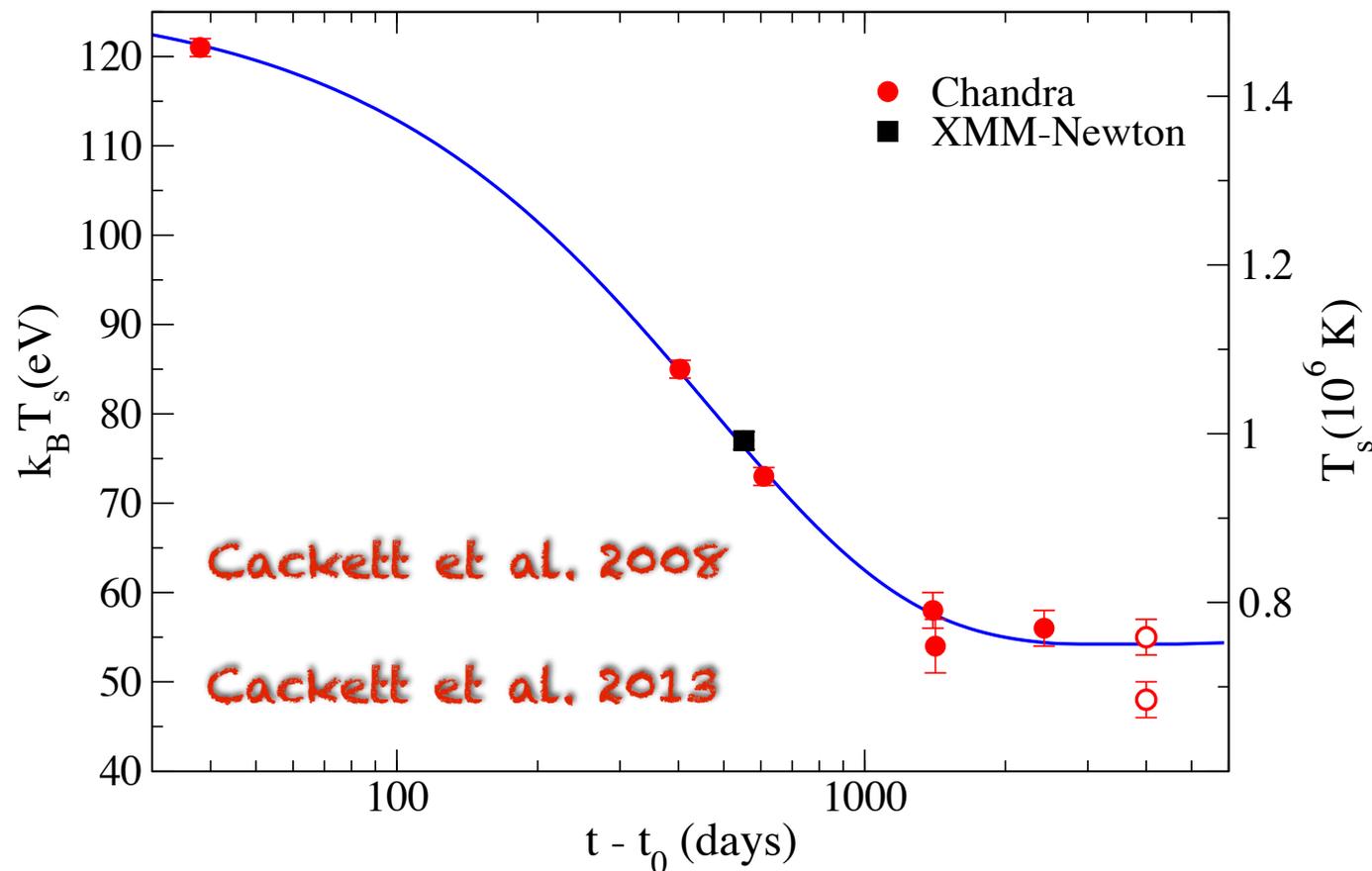
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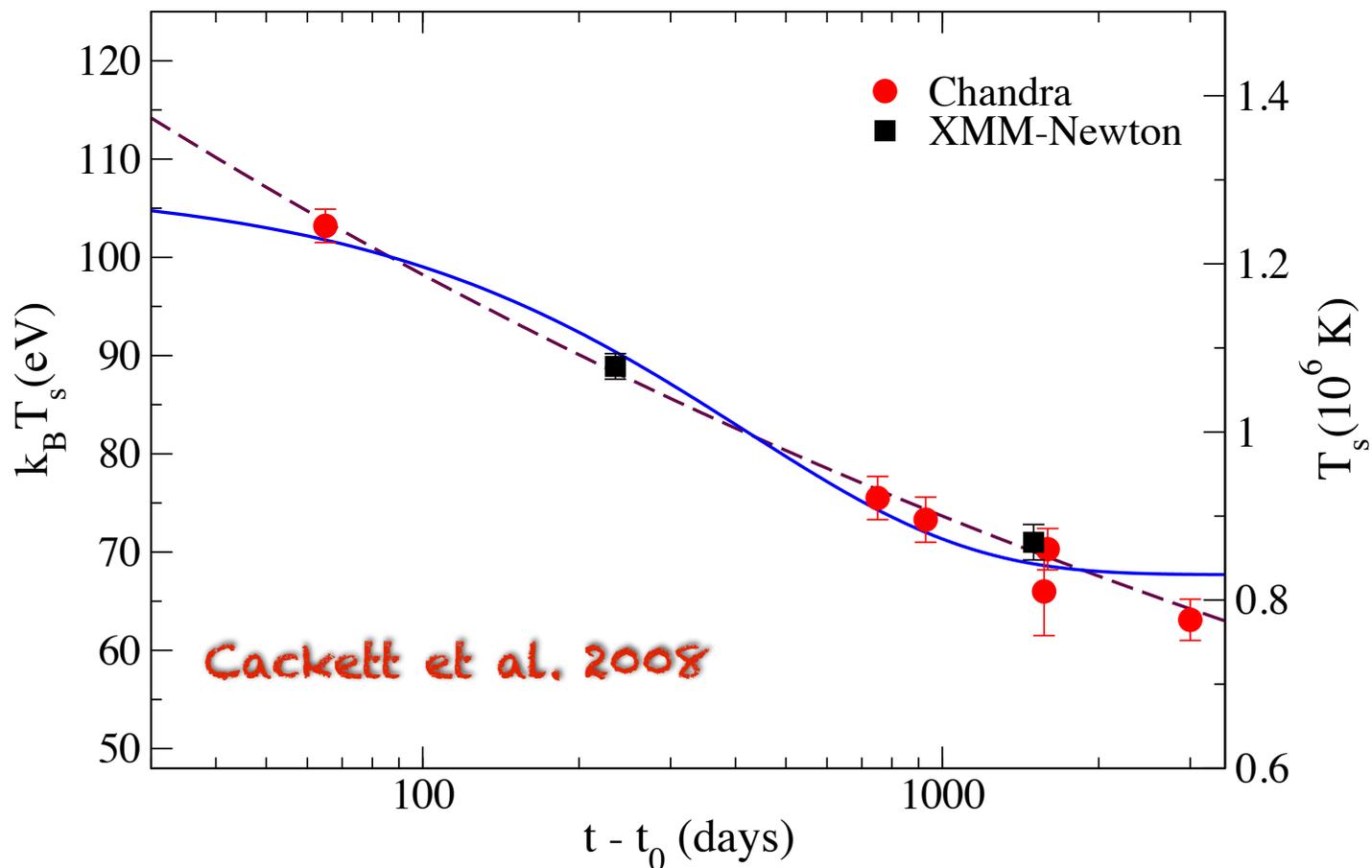


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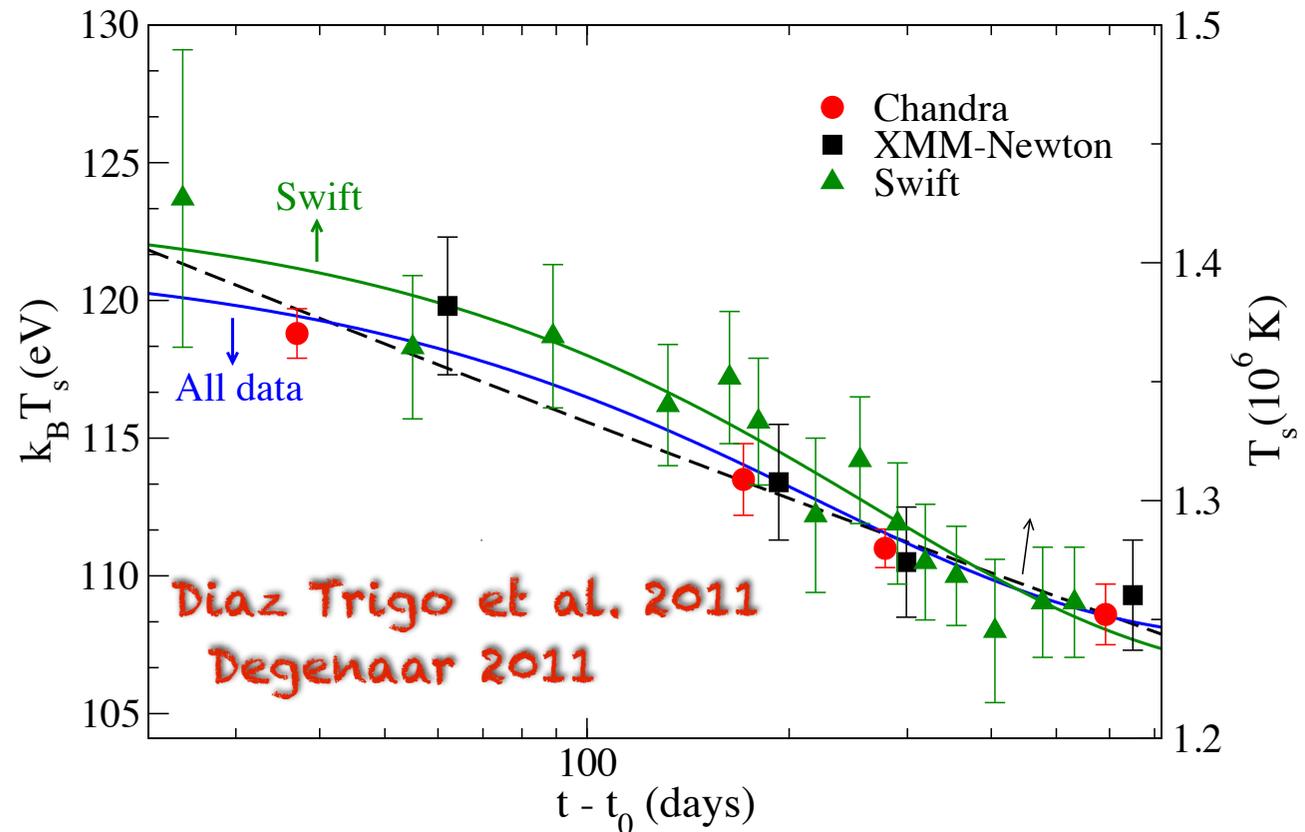


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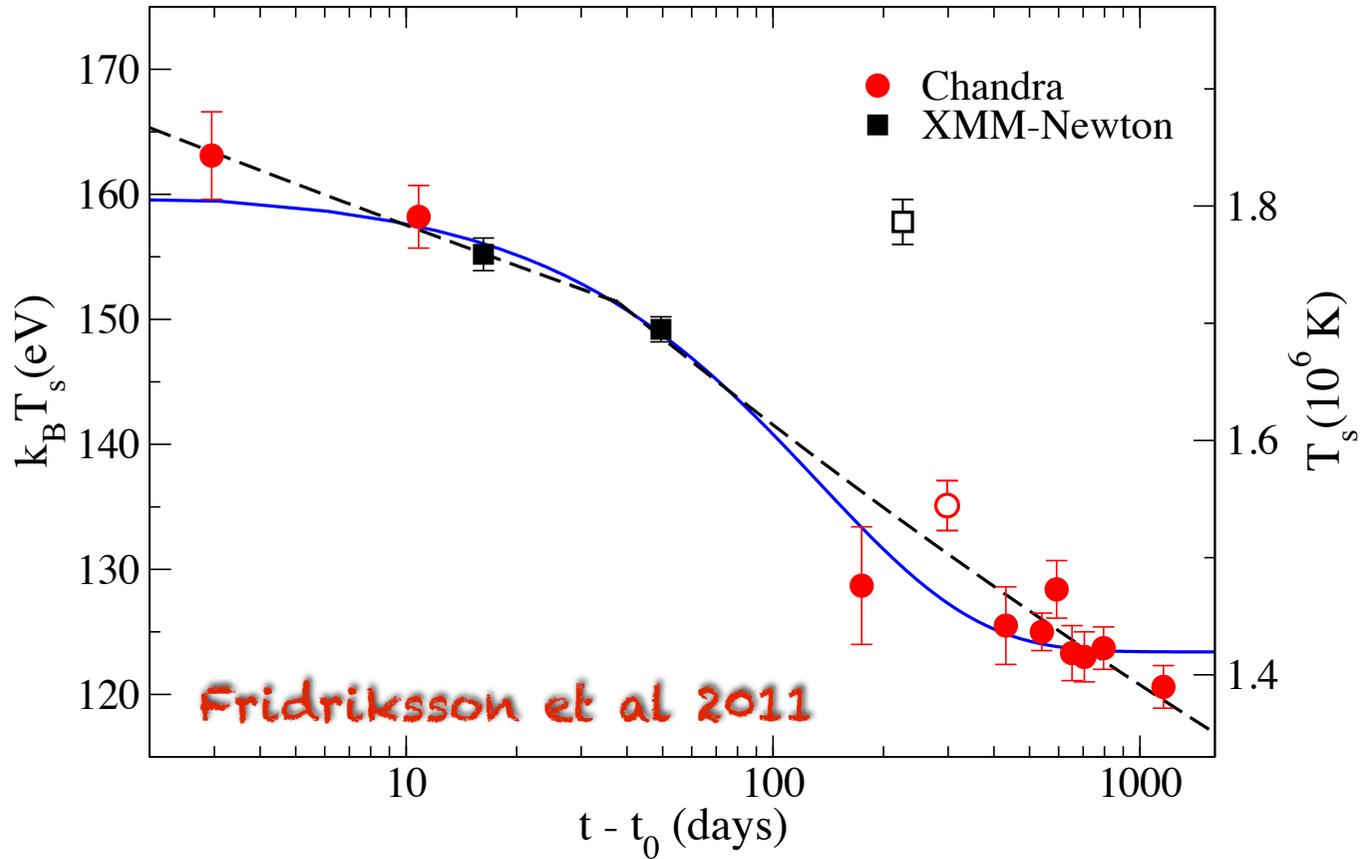


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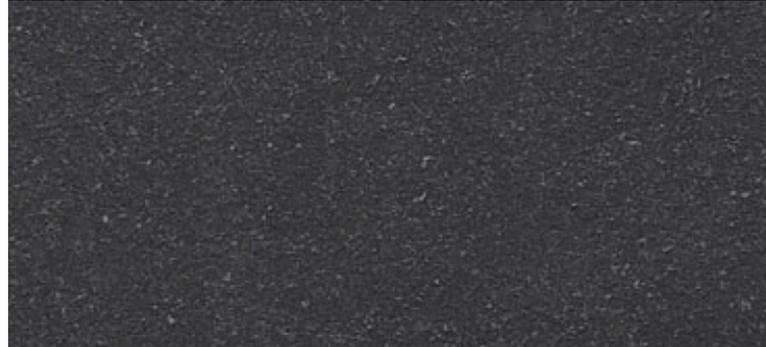
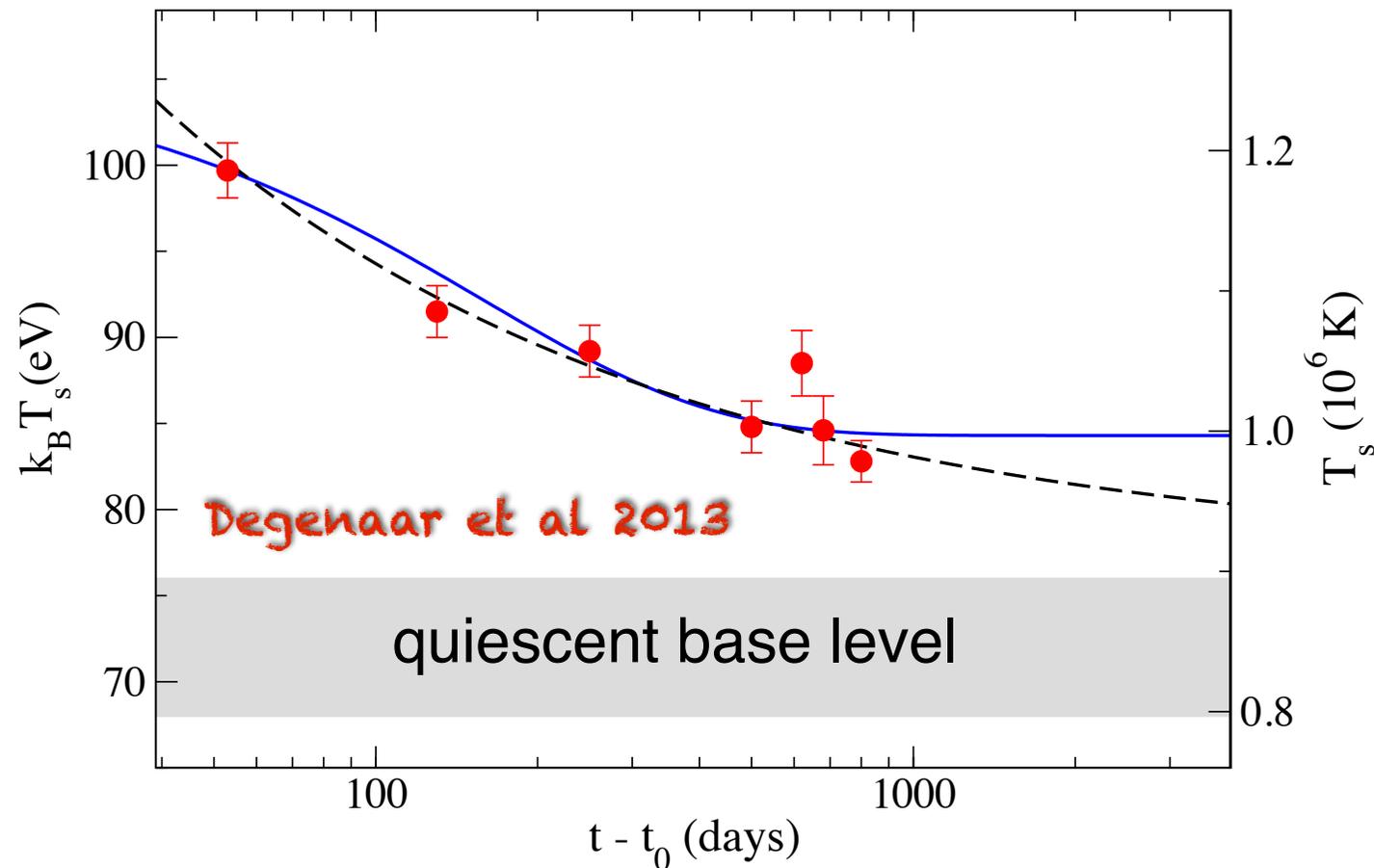
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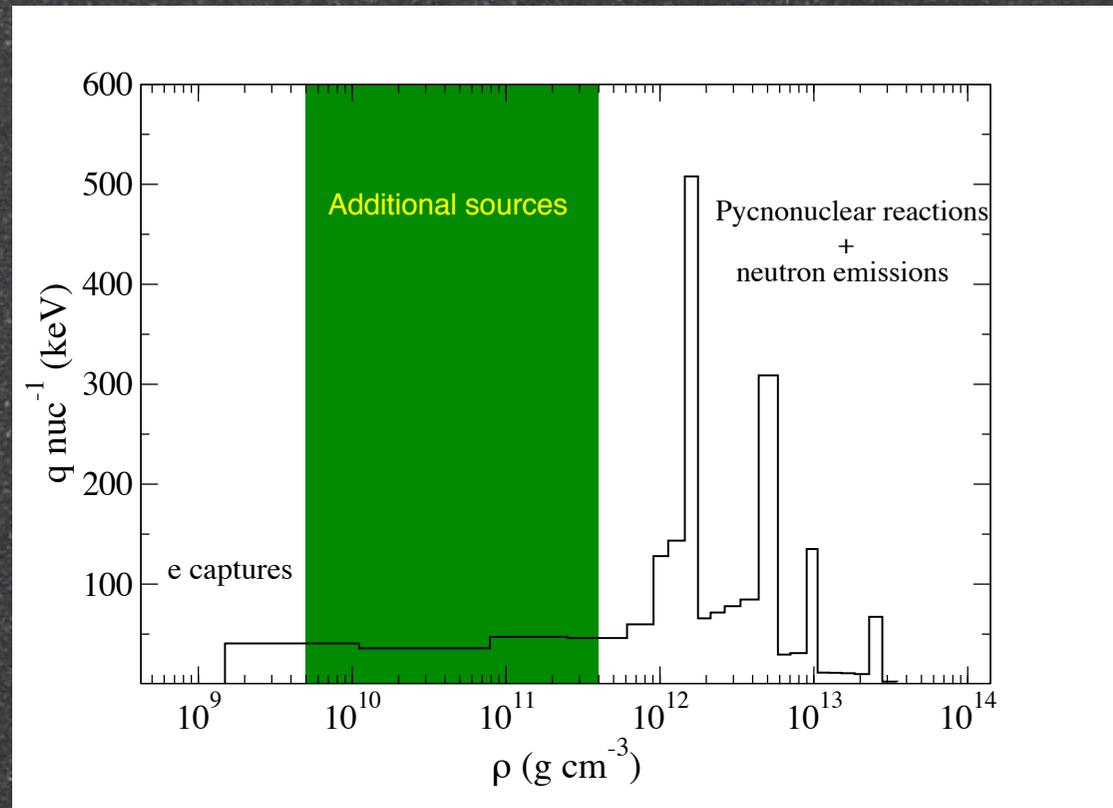


# Modeling the crust thermal evolution

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## Heat sources

$$Q_{\dot{m}} = \dot{m} \int_{\rho_s}^{\rho} q(\rho) d\rho = \dot{m} \sum_j q(\rho_j)$$



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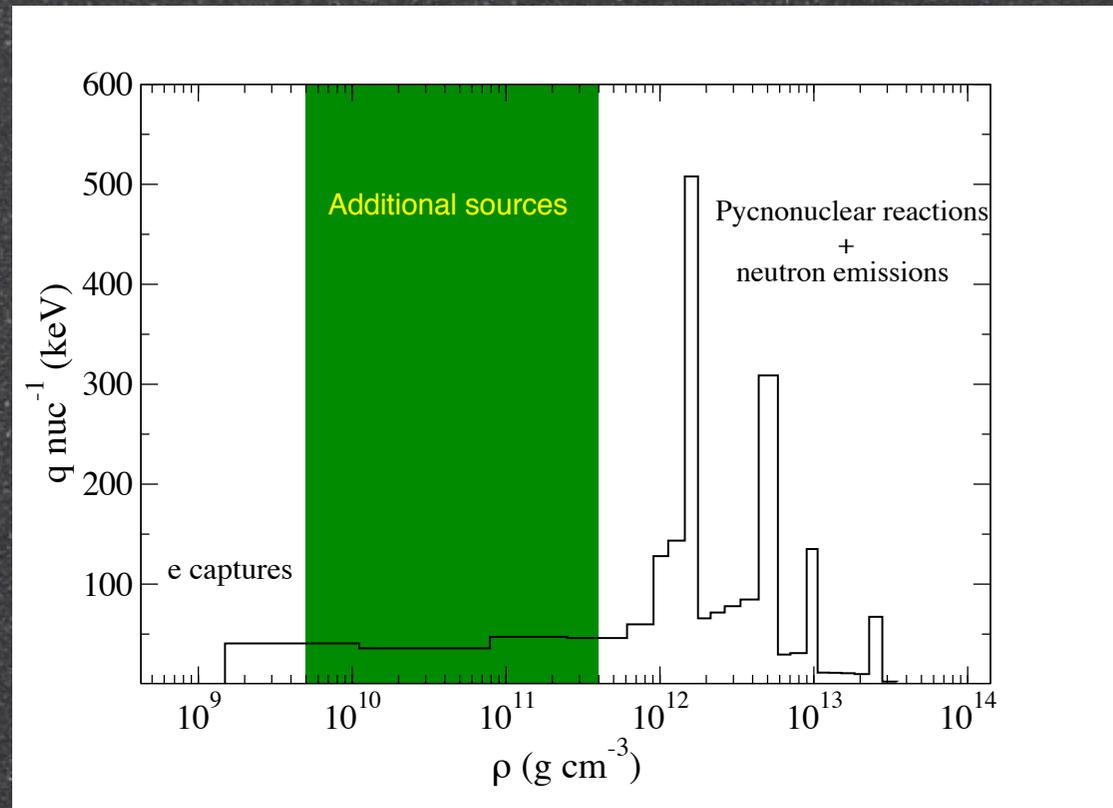
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NS model includes: realistic accreted EoS, neutrino emissivity, specific heat, thermal conductivity, etc.

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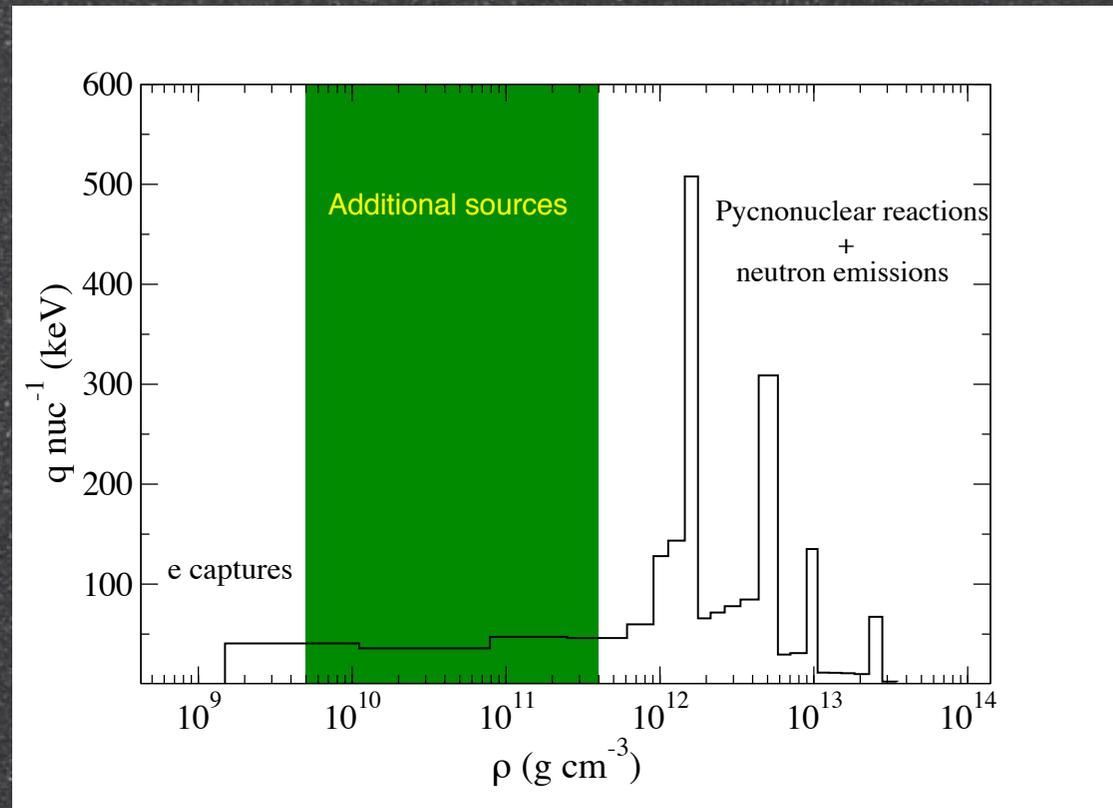
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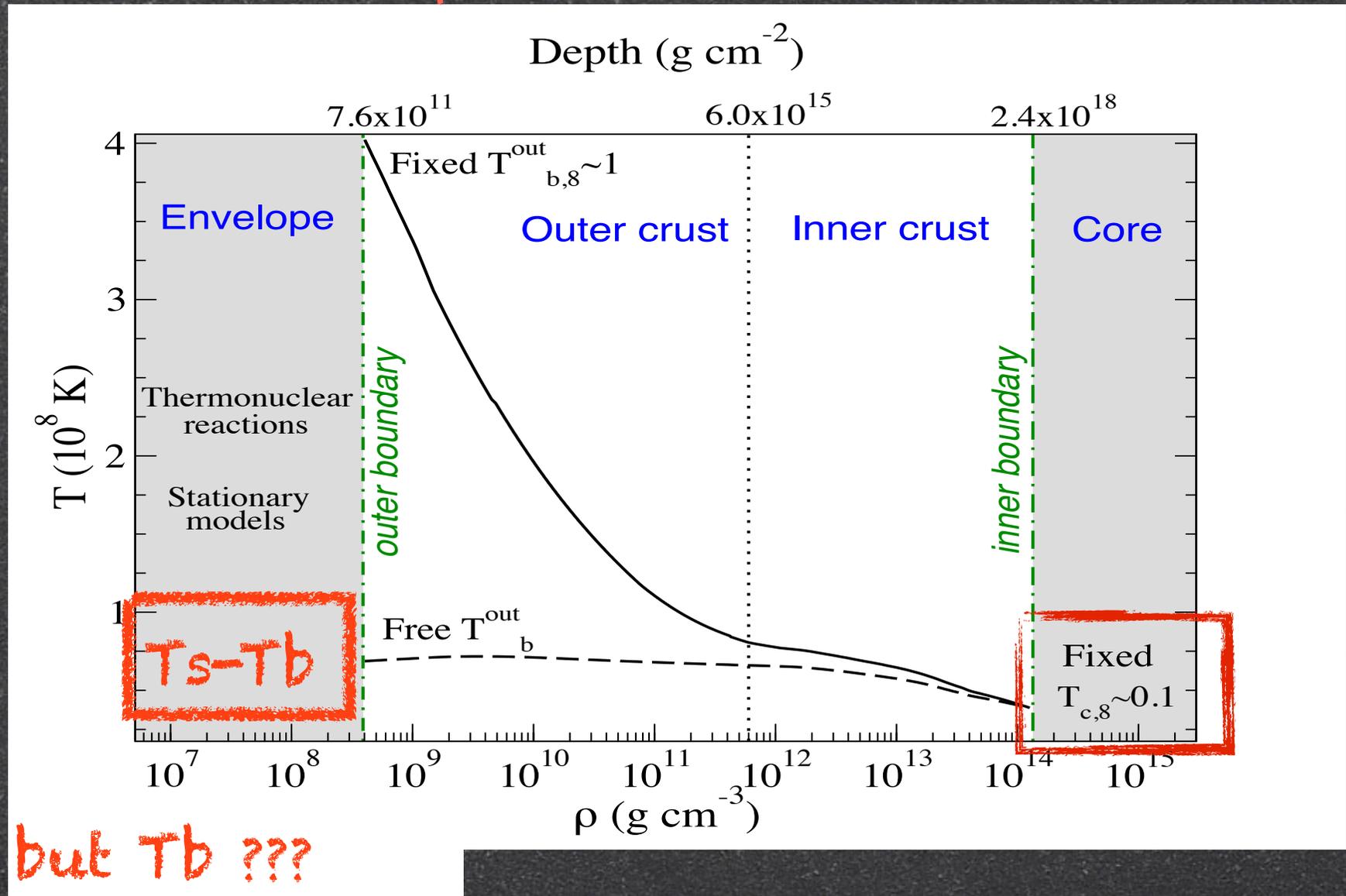


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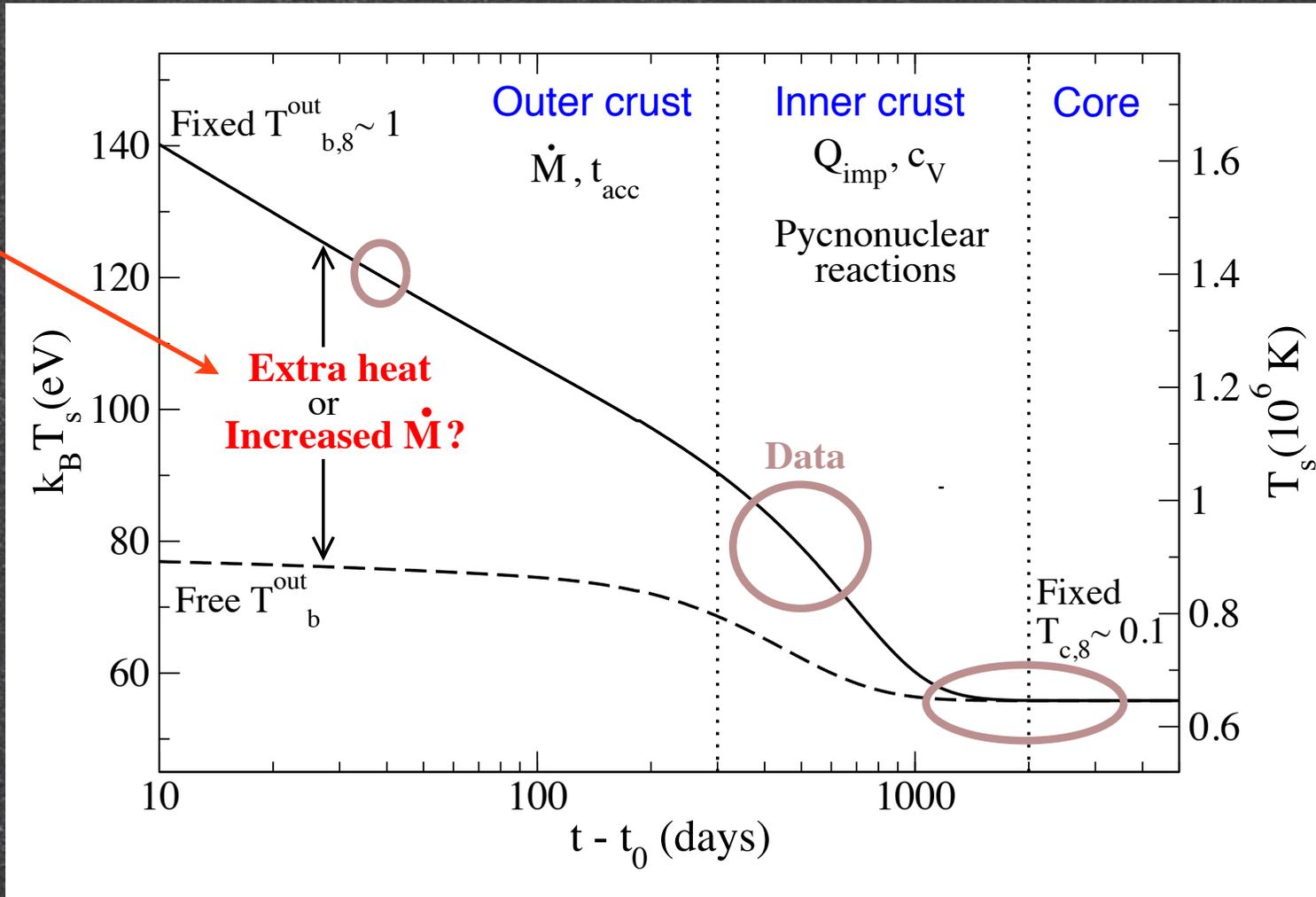
BC { inner  $T_{\text{core}}$  constant  
outer  $T_{\text{envelope}}??$

# MXB 1659-29: Textbook crustal cooler

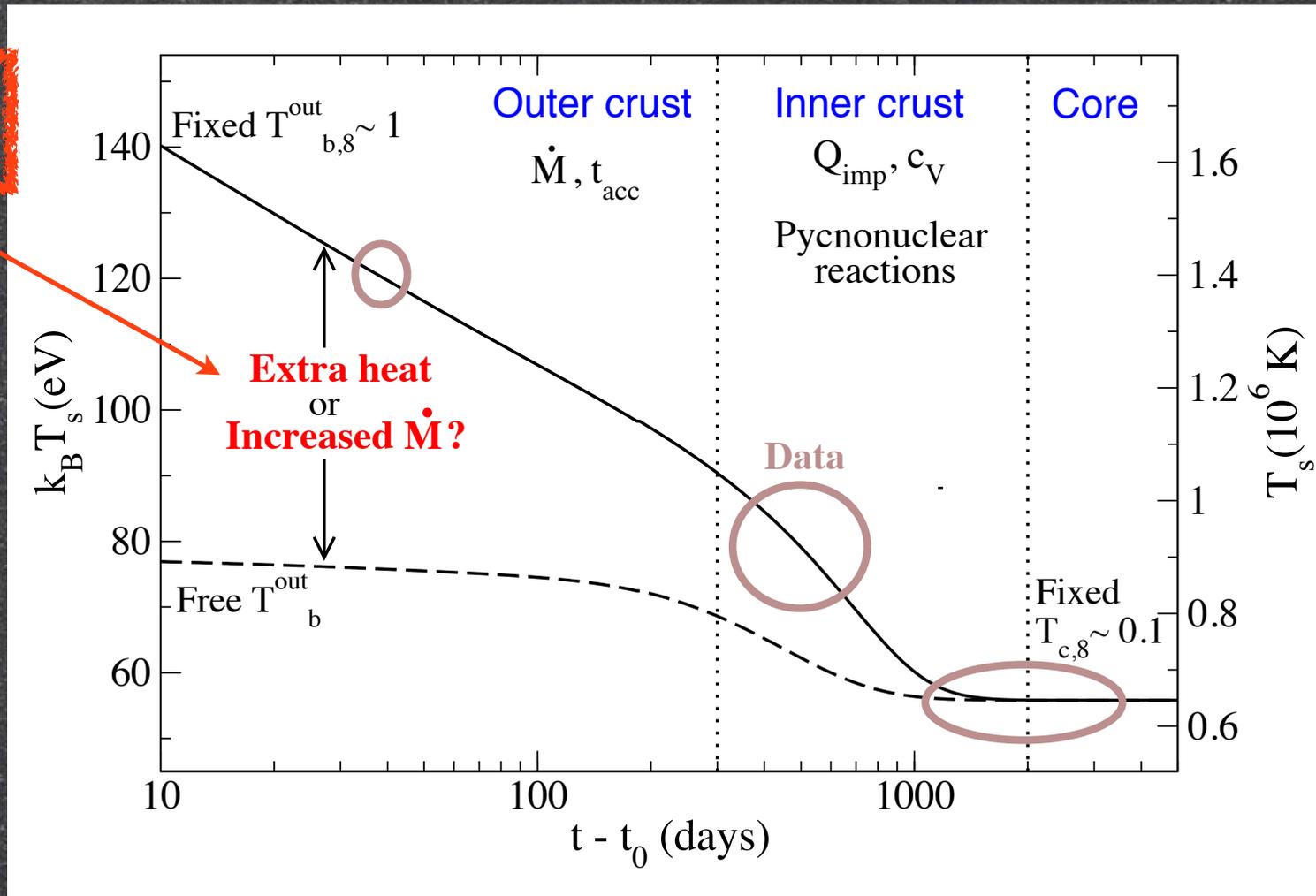
Initial thermal profile



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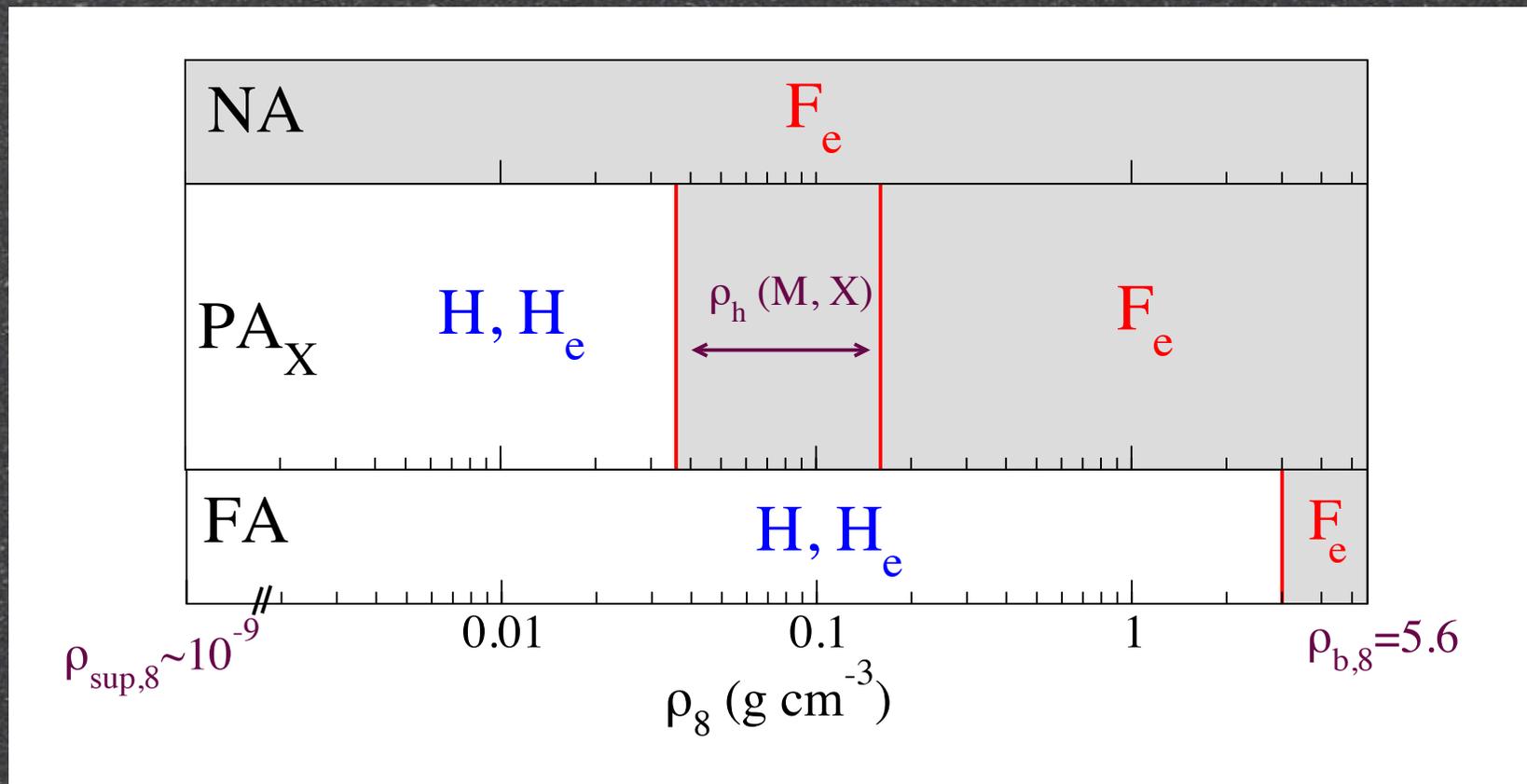




# The envelope

in quiescence

in accretion



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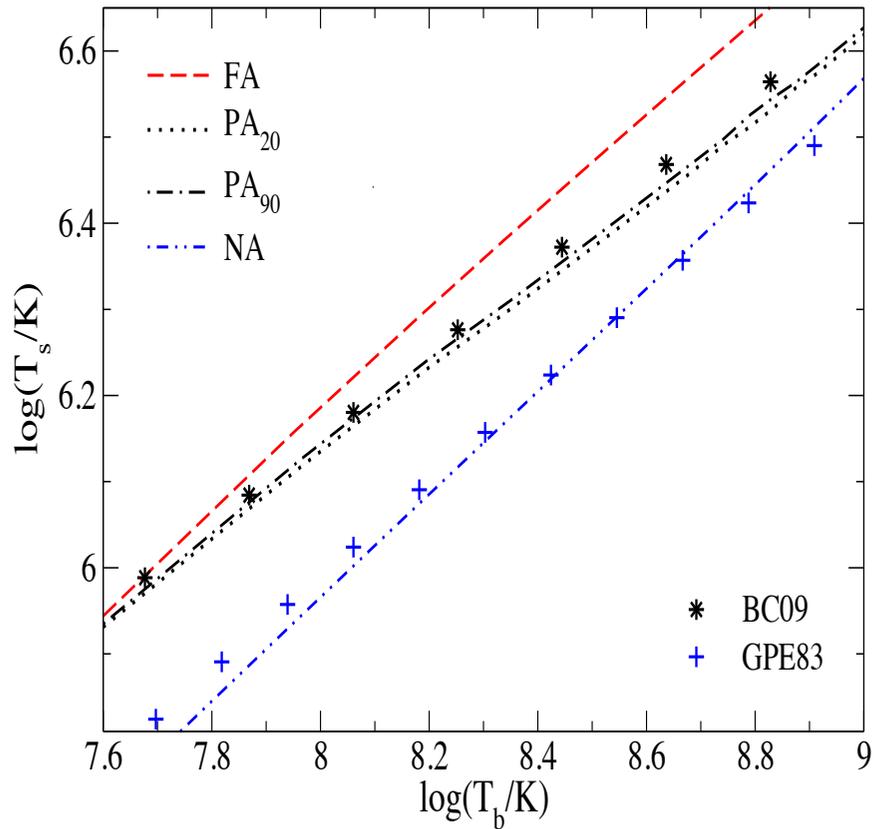
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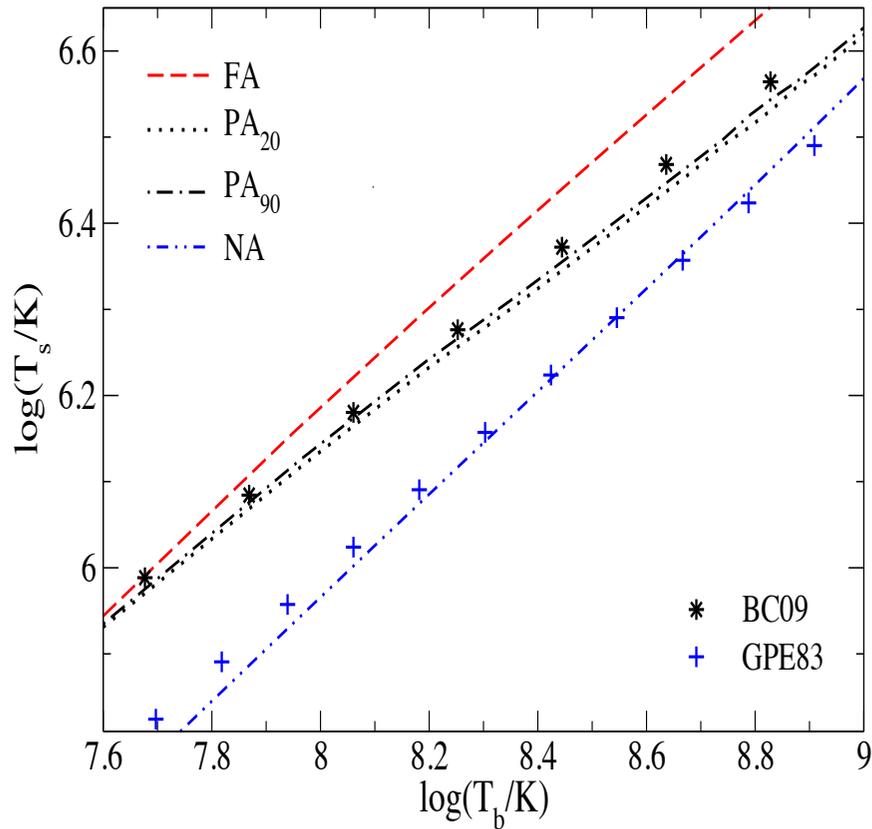
in accretion



stationary solutions

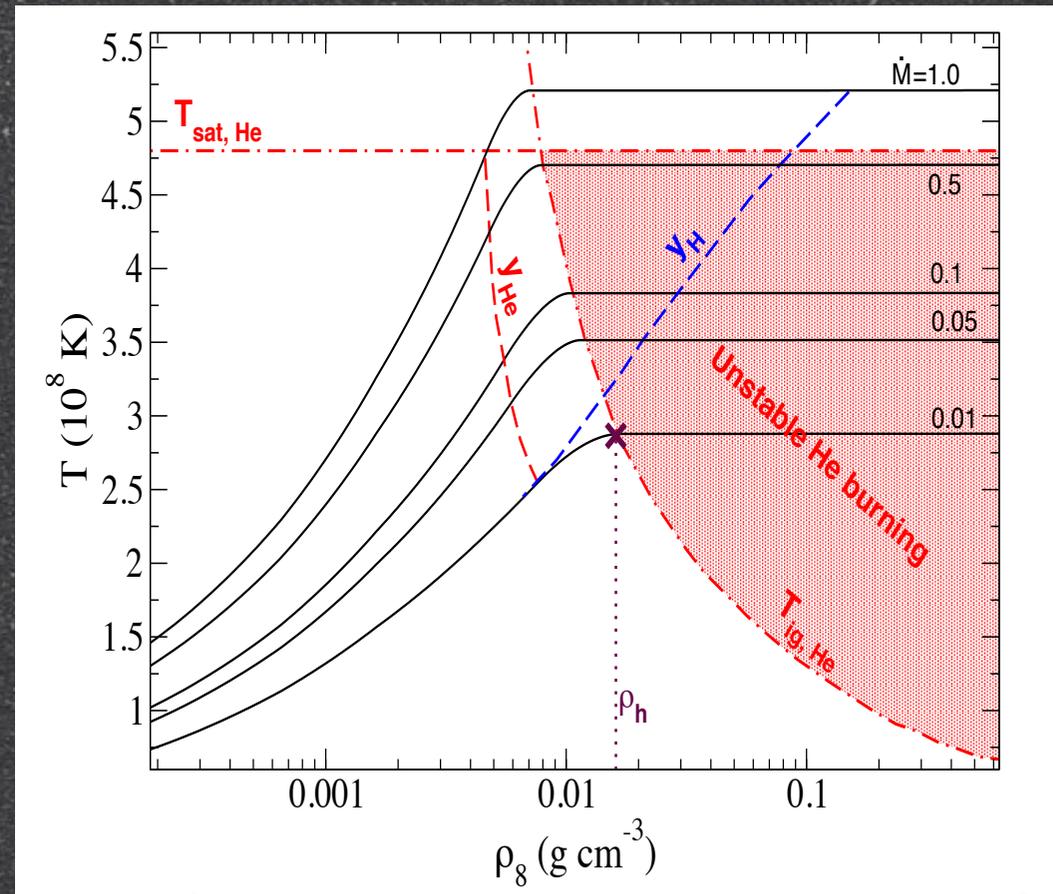
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in quiescence



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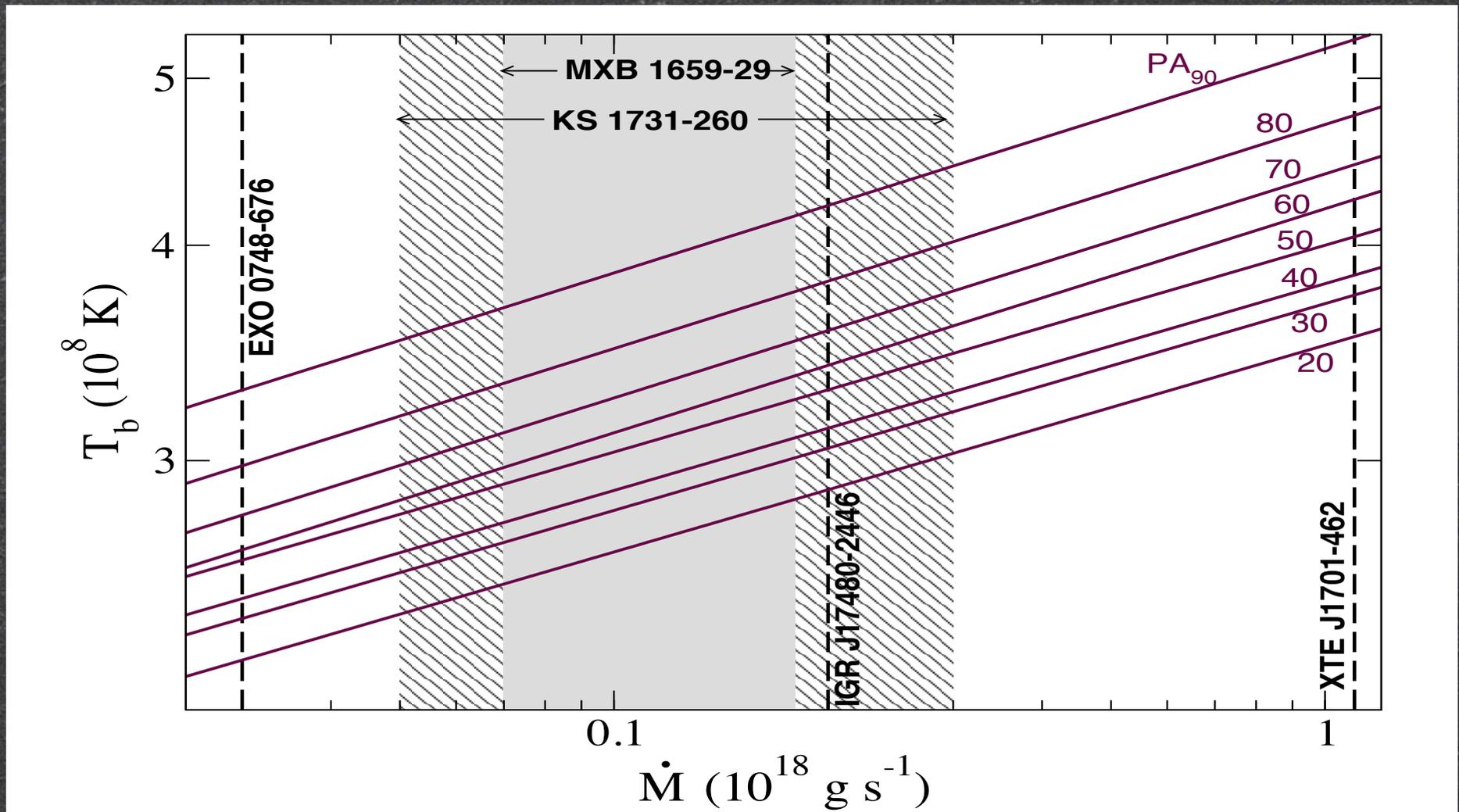
in accretion



including hot CNO-cycle and  
3-alpha generation rates  
Bildsten (97)

# Relations $T_b(\dot{M})$

$T_b$  is modified in the active phase through  $\dot{M}$   
→ stationary profiles for  $T_b(\dot{M})$

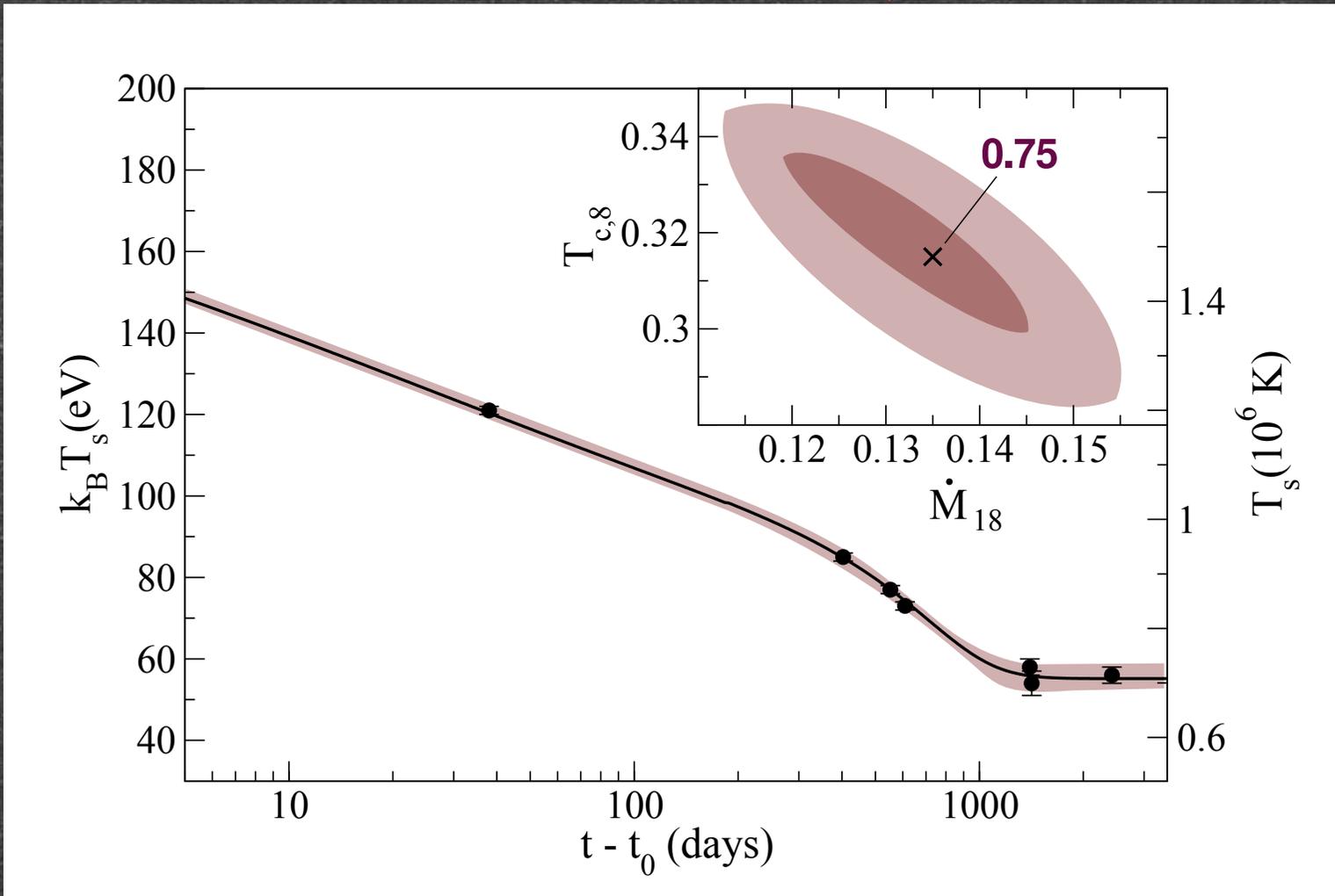


"Crustal coolers"

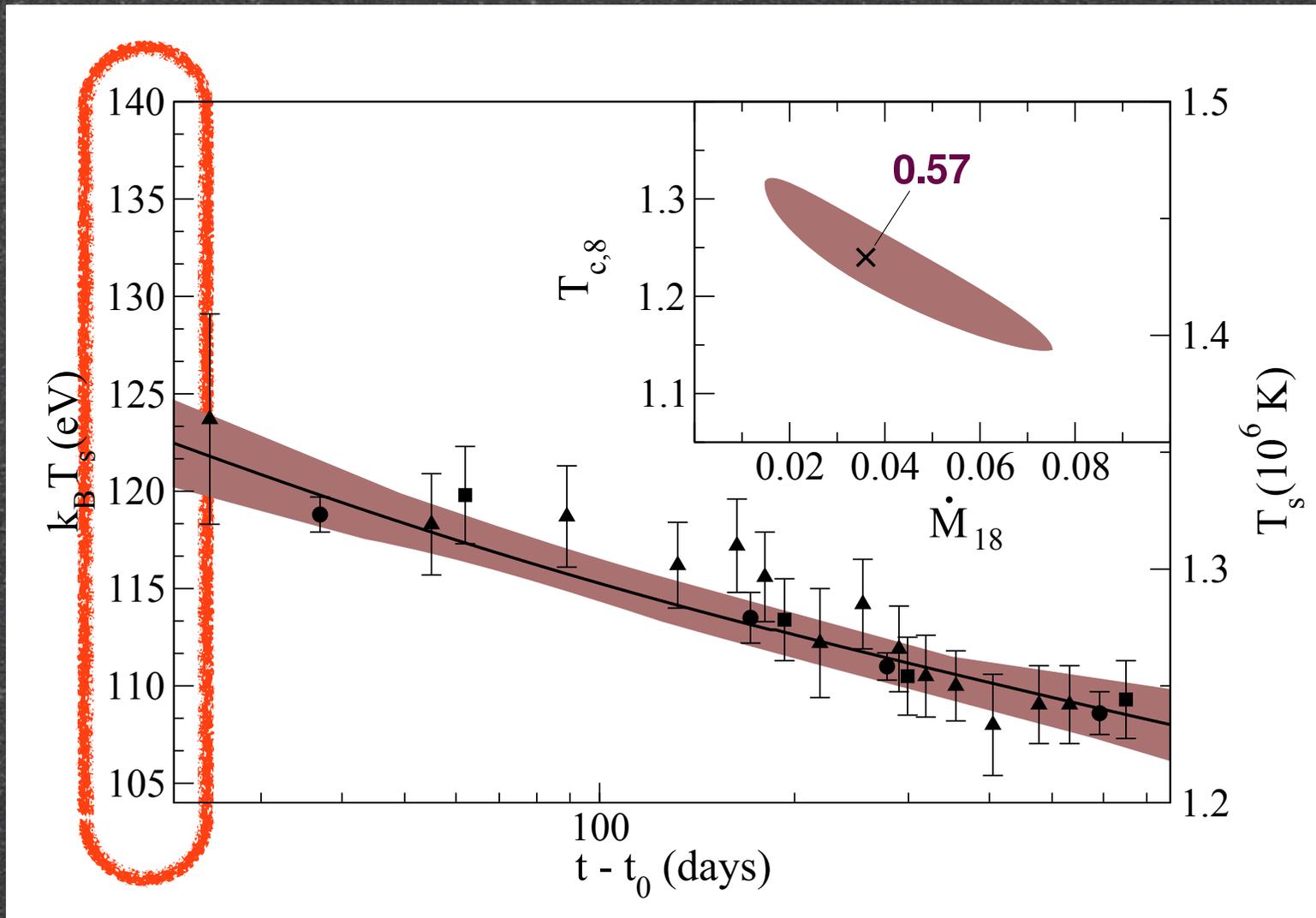
# MXB 1659-29: Textbook crystal cooler

Tb( $\dot{M}$ )

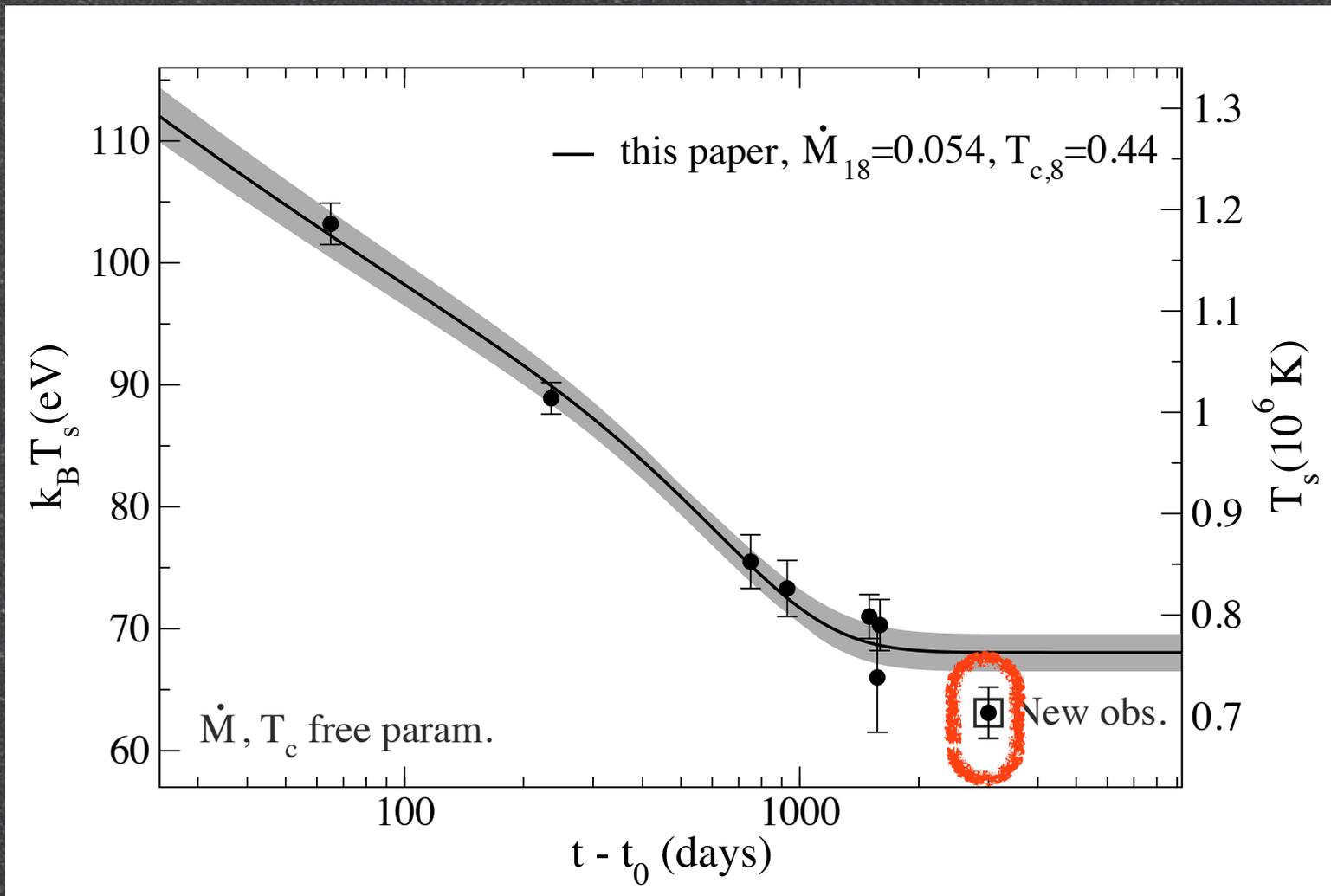
→ Reduces the # of free parameters



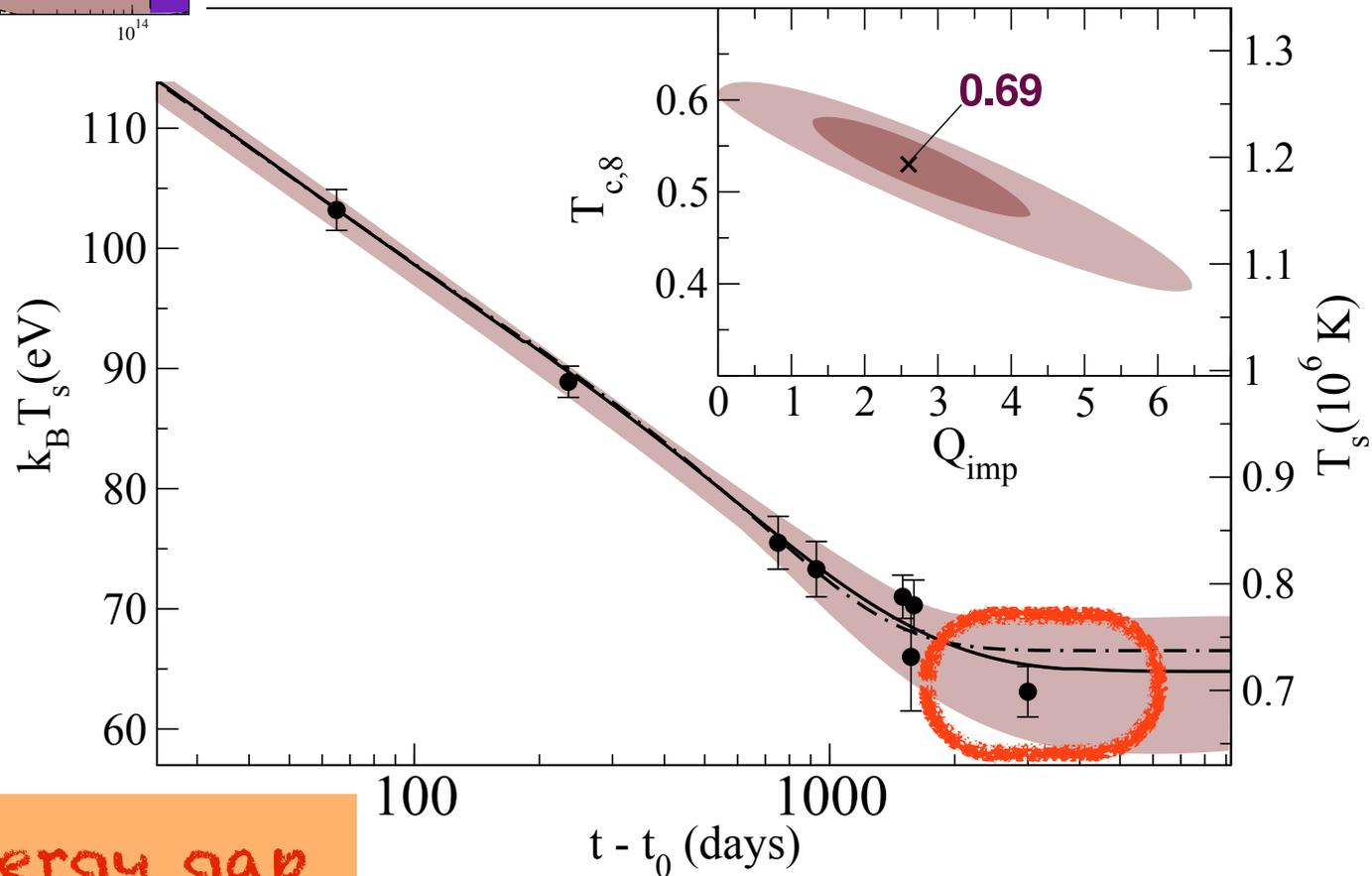
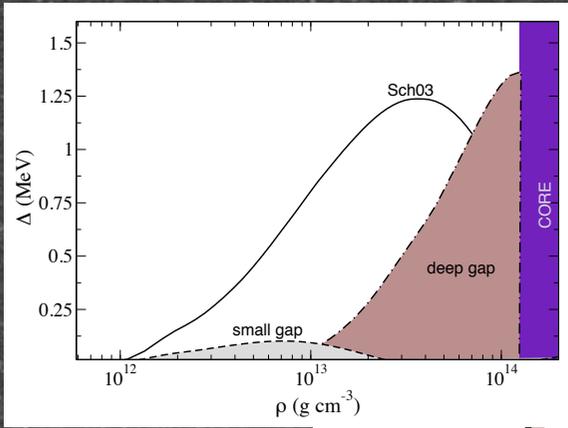
# EXO 0748-676: Hot crustal cooler



But...KS 1731-260  
still cooling?



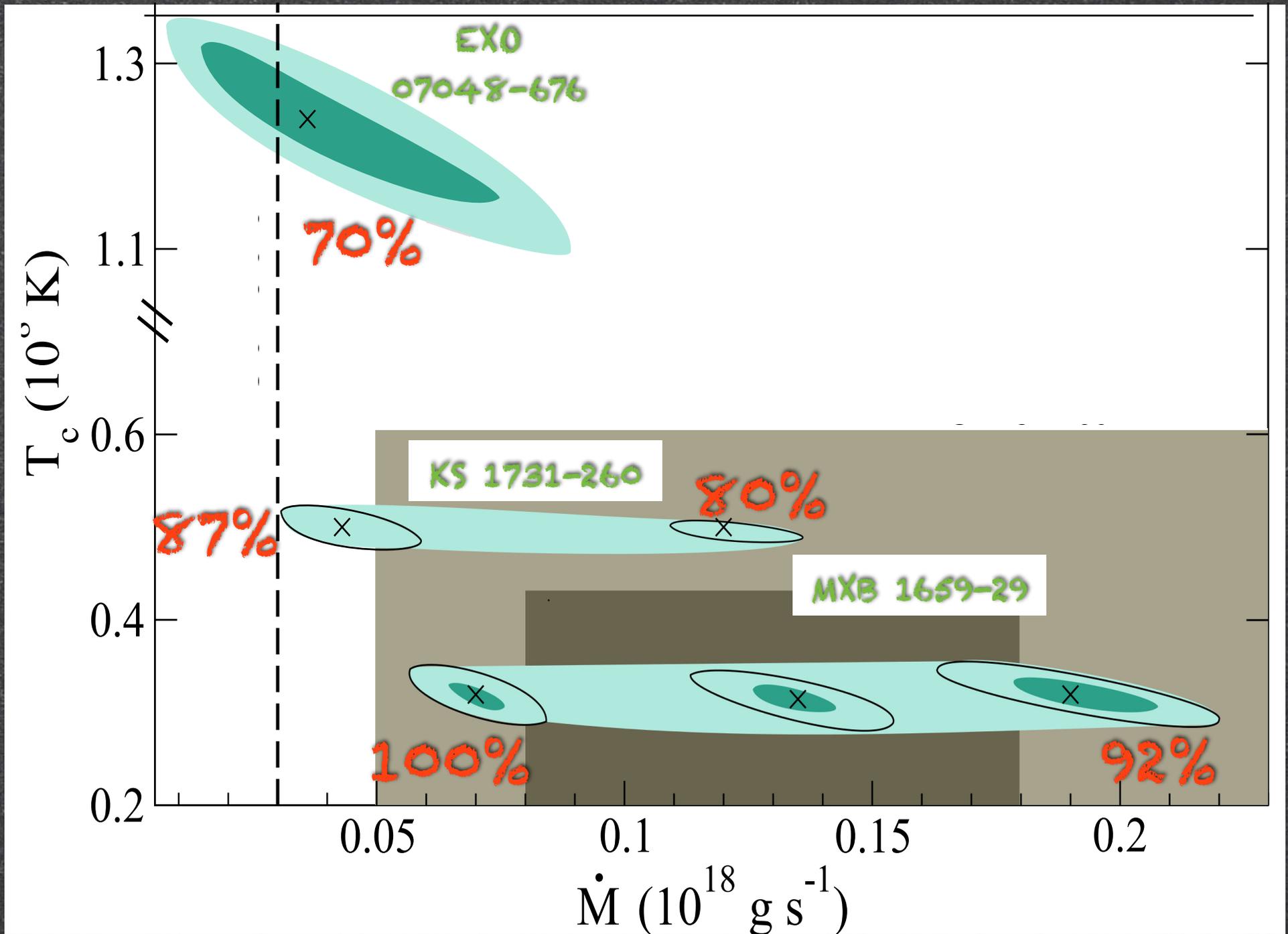
# KS 1731-260 .... adjusting model



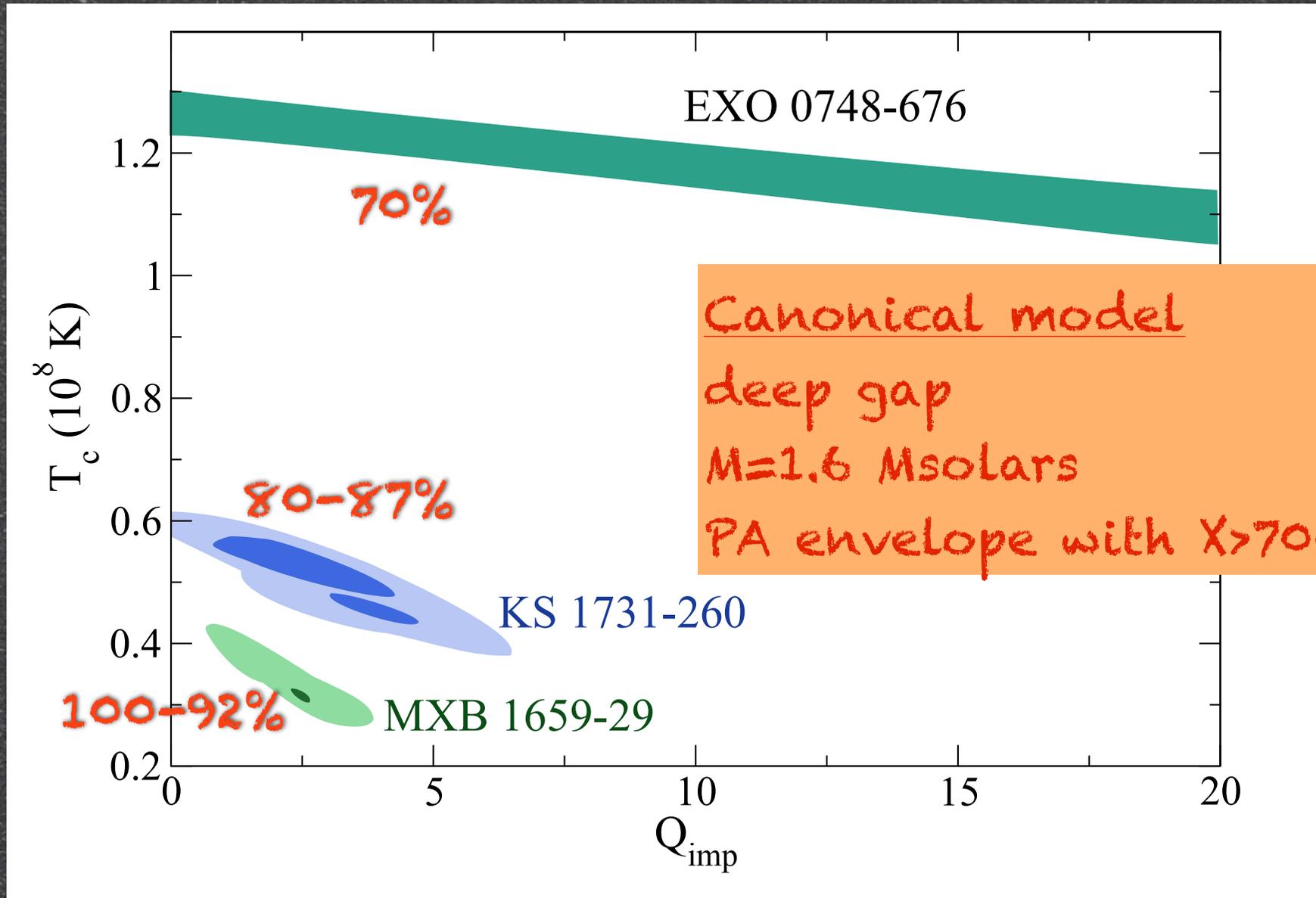
"deep" energy gap  
for neutron SF!!

Longer relaxation time

# Constraints to crust/envelope

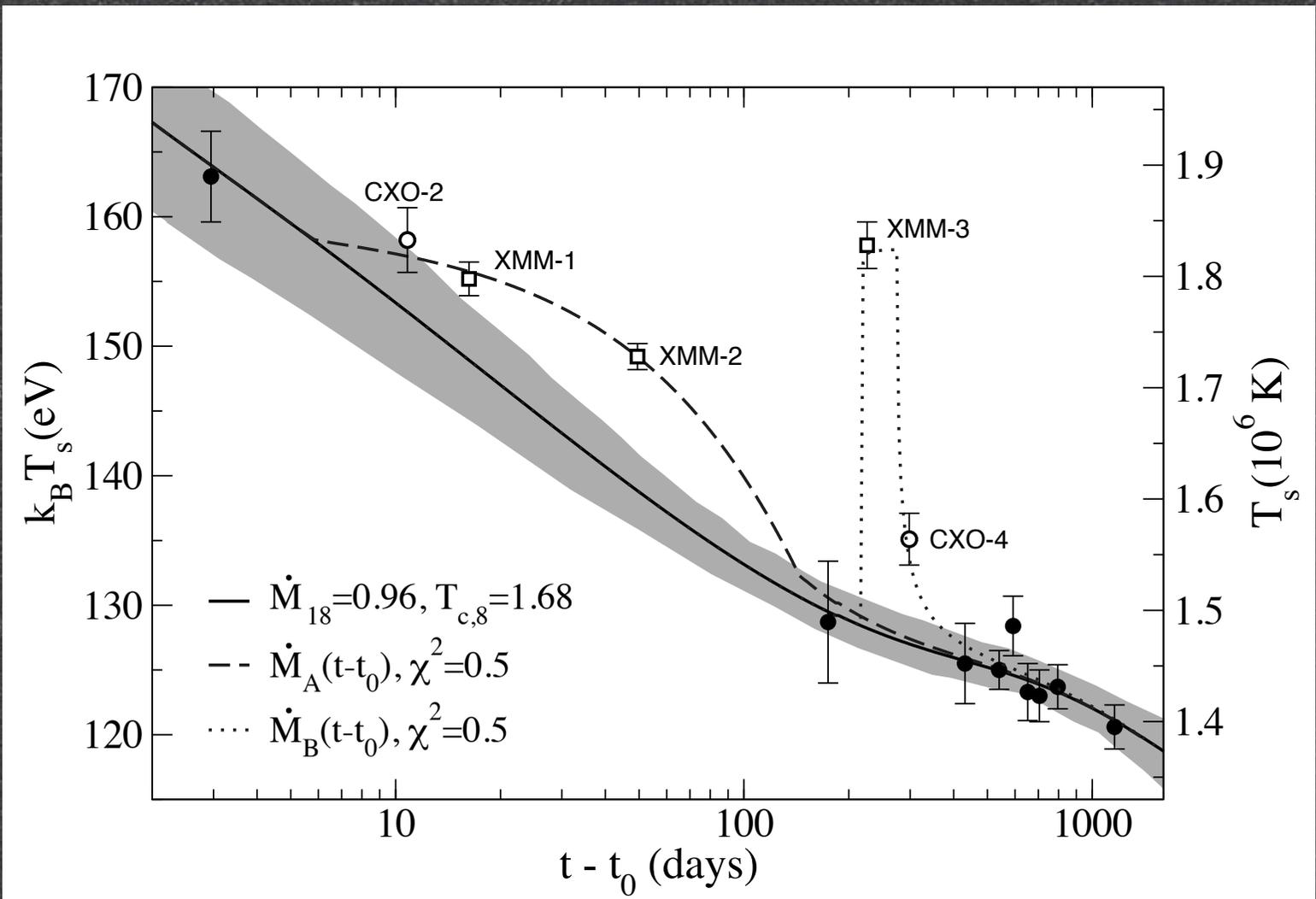


# Crustal coolers



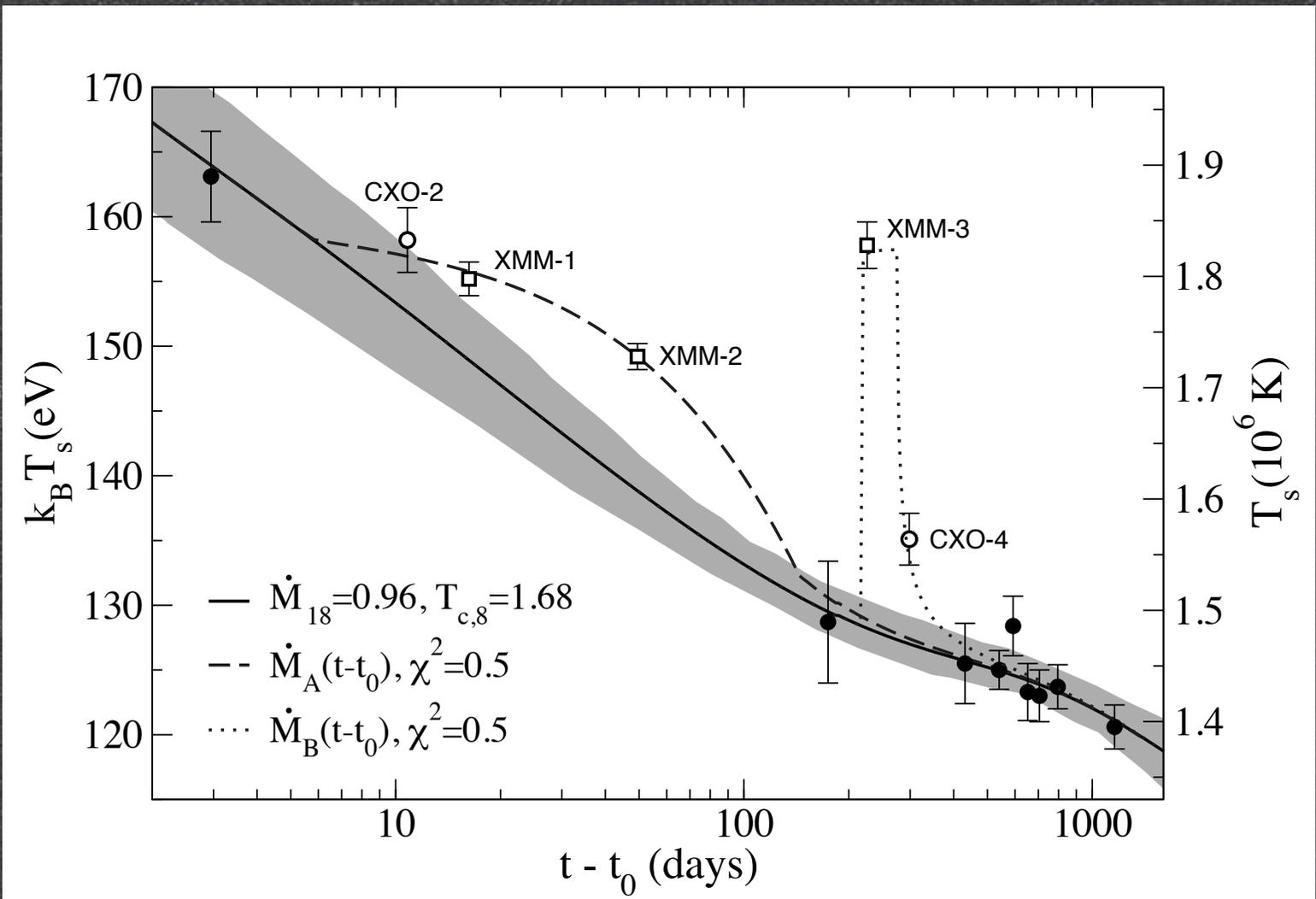
"Beyond" crustal  
cooling

# XTE 1701-462: VARIABILITY



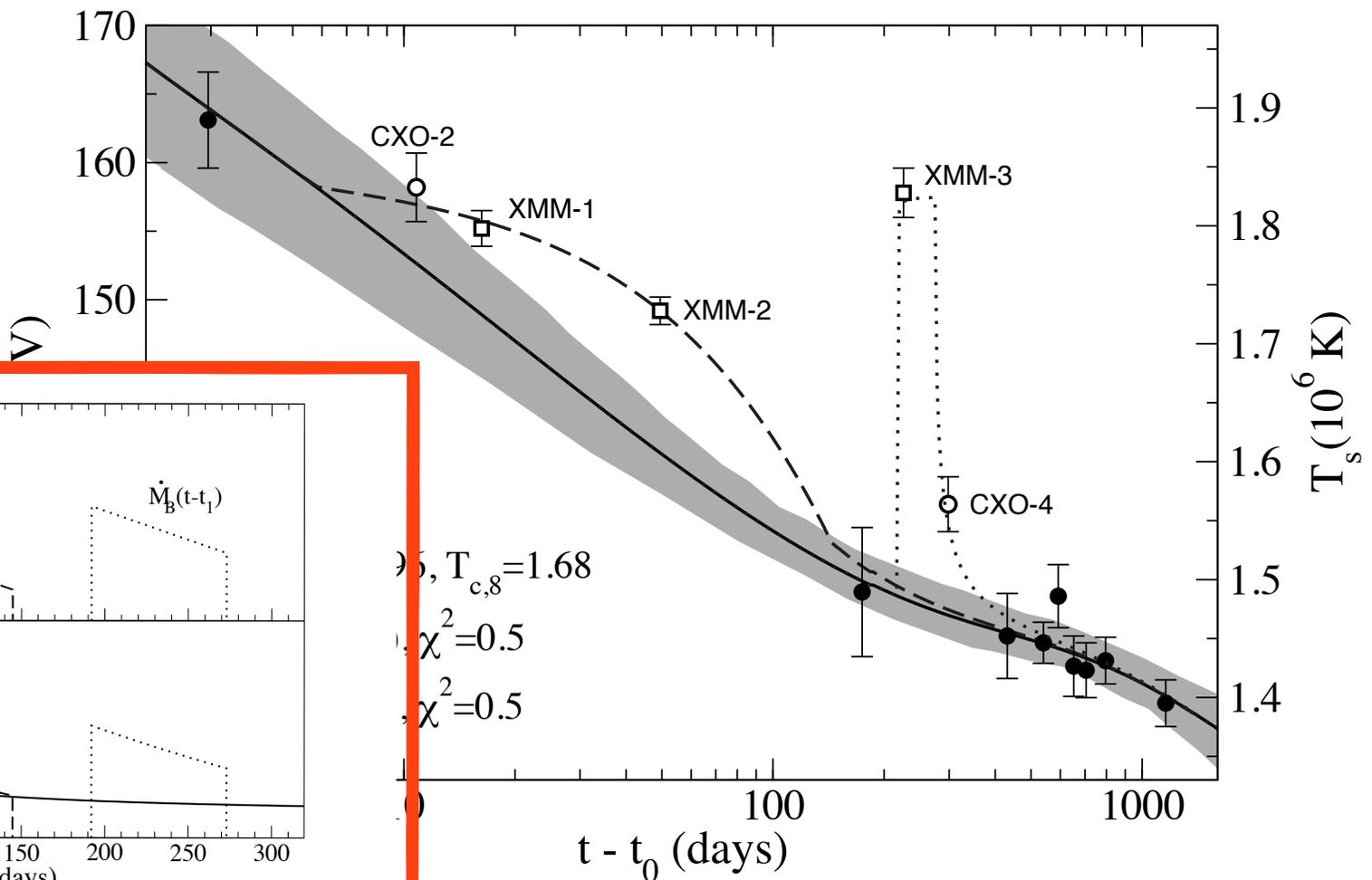
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residual accretion?

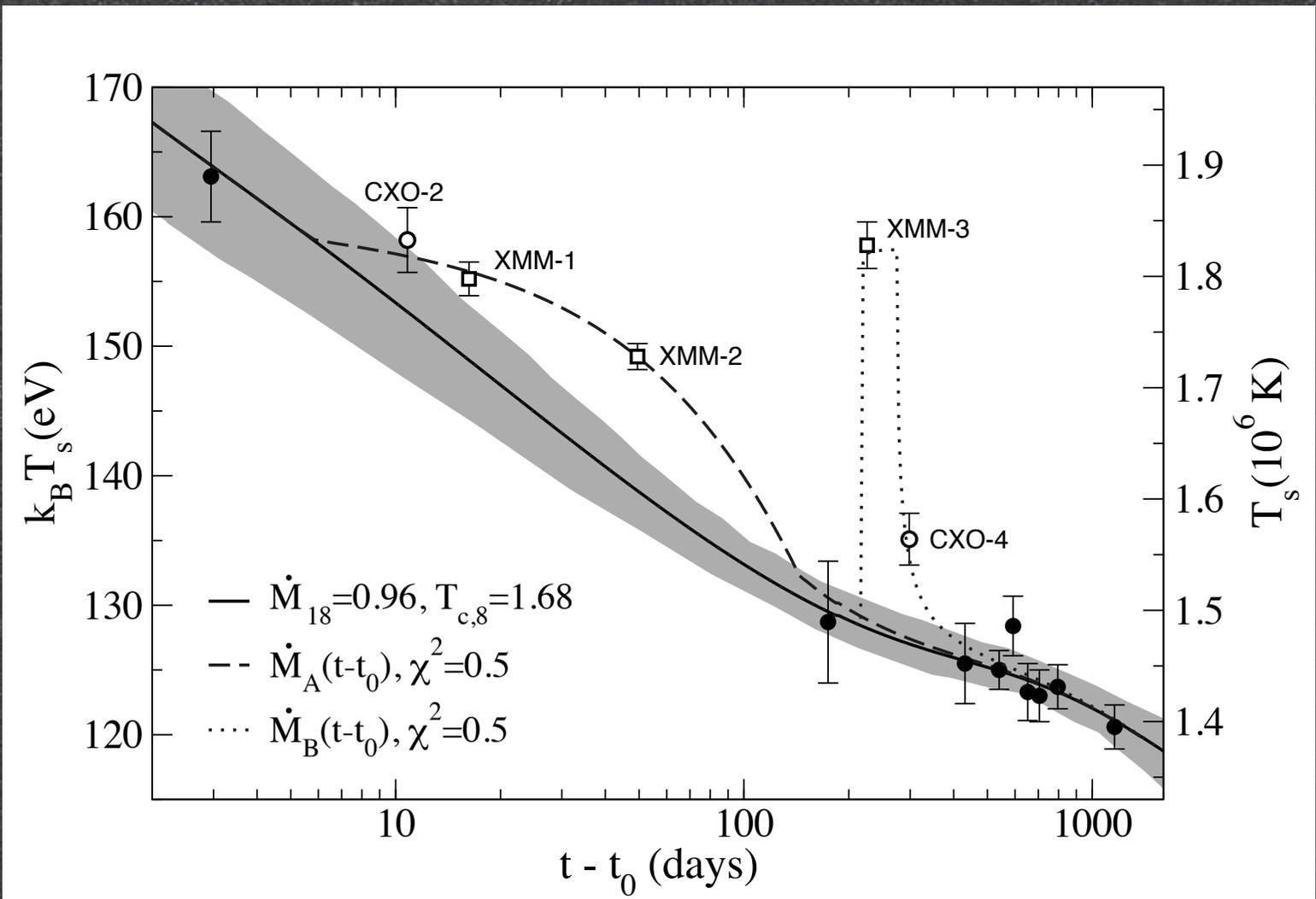


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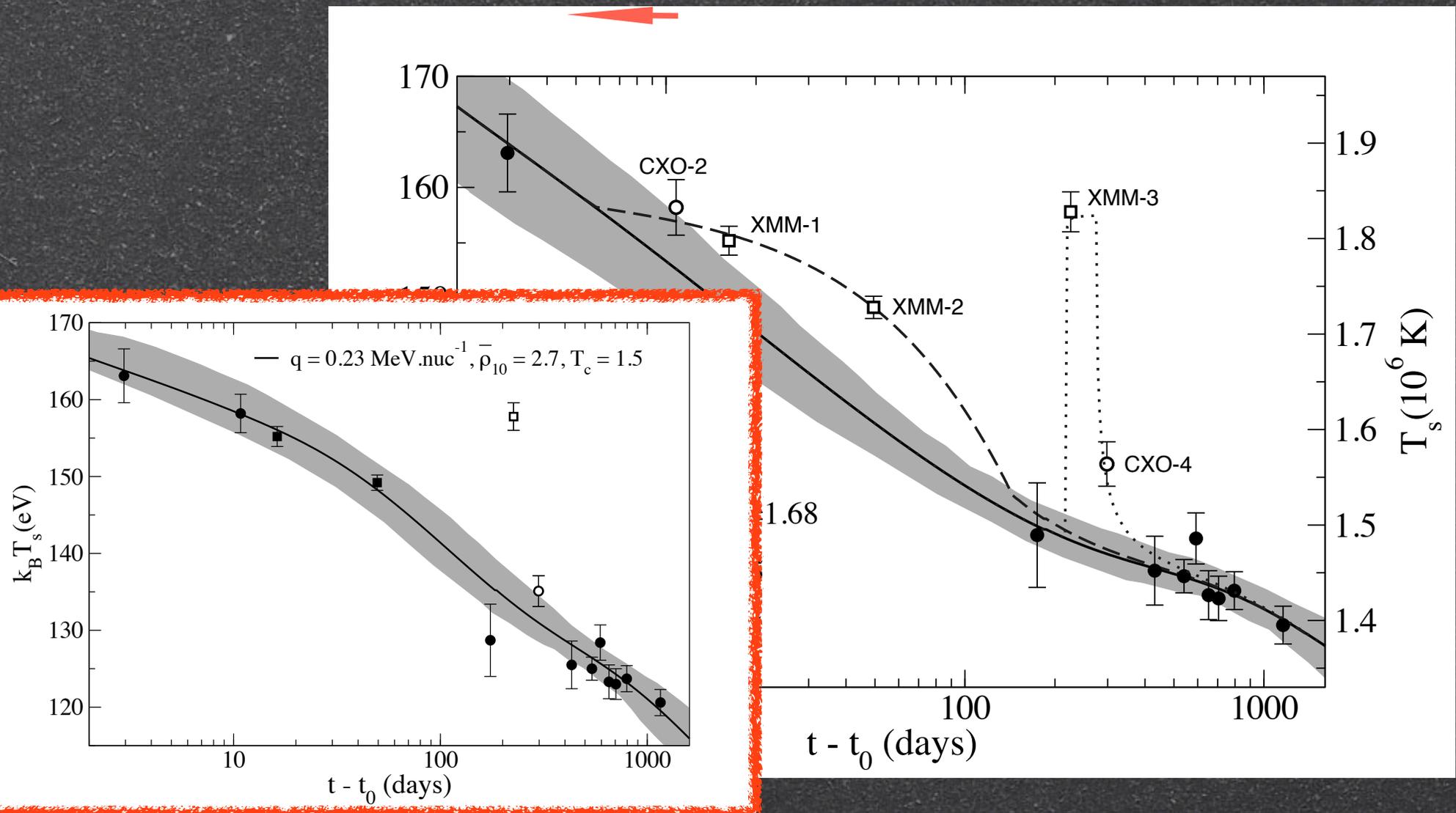


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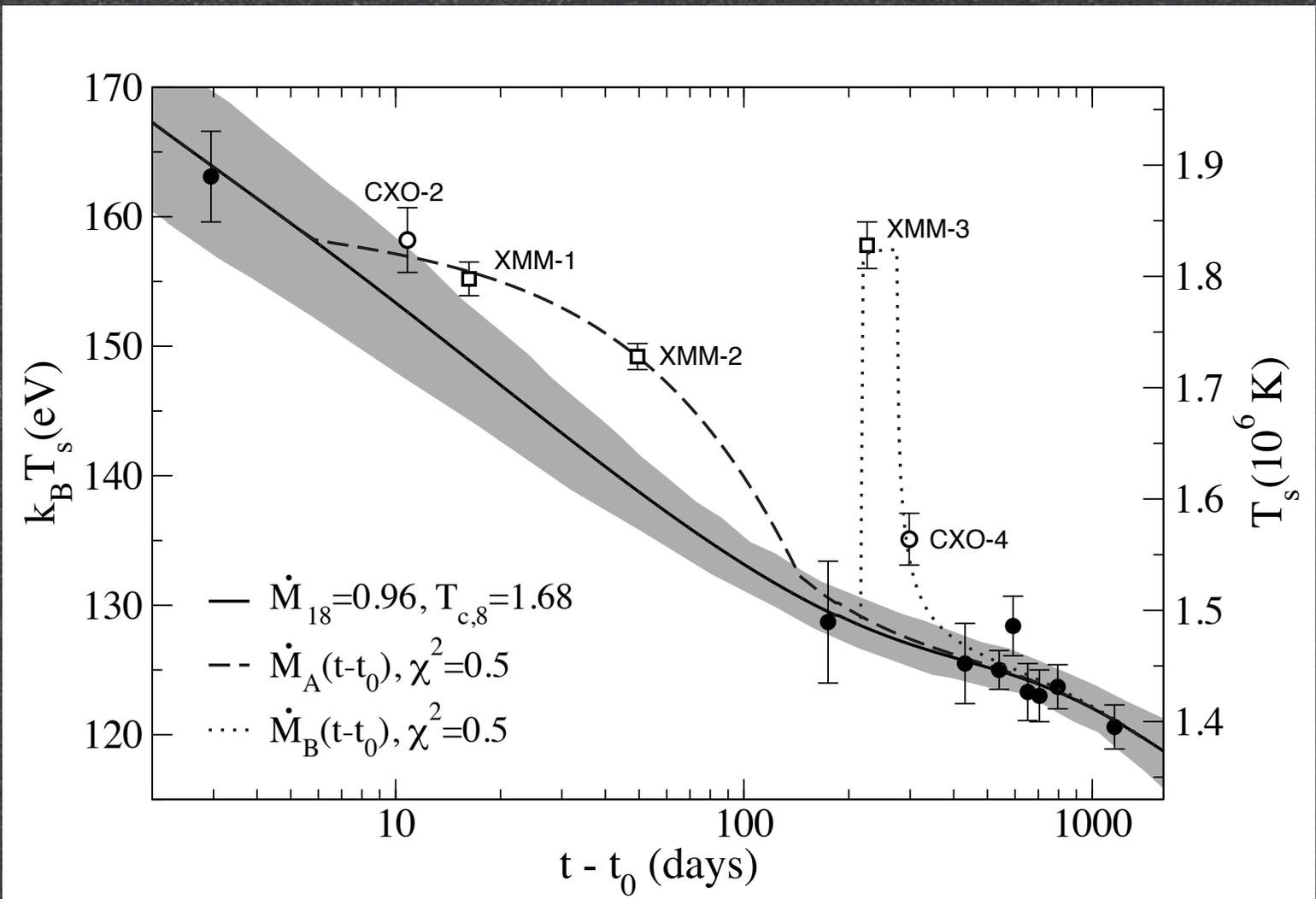


# XTE 1701-462: VARIABILITY

other heat sources @ shallow depths?

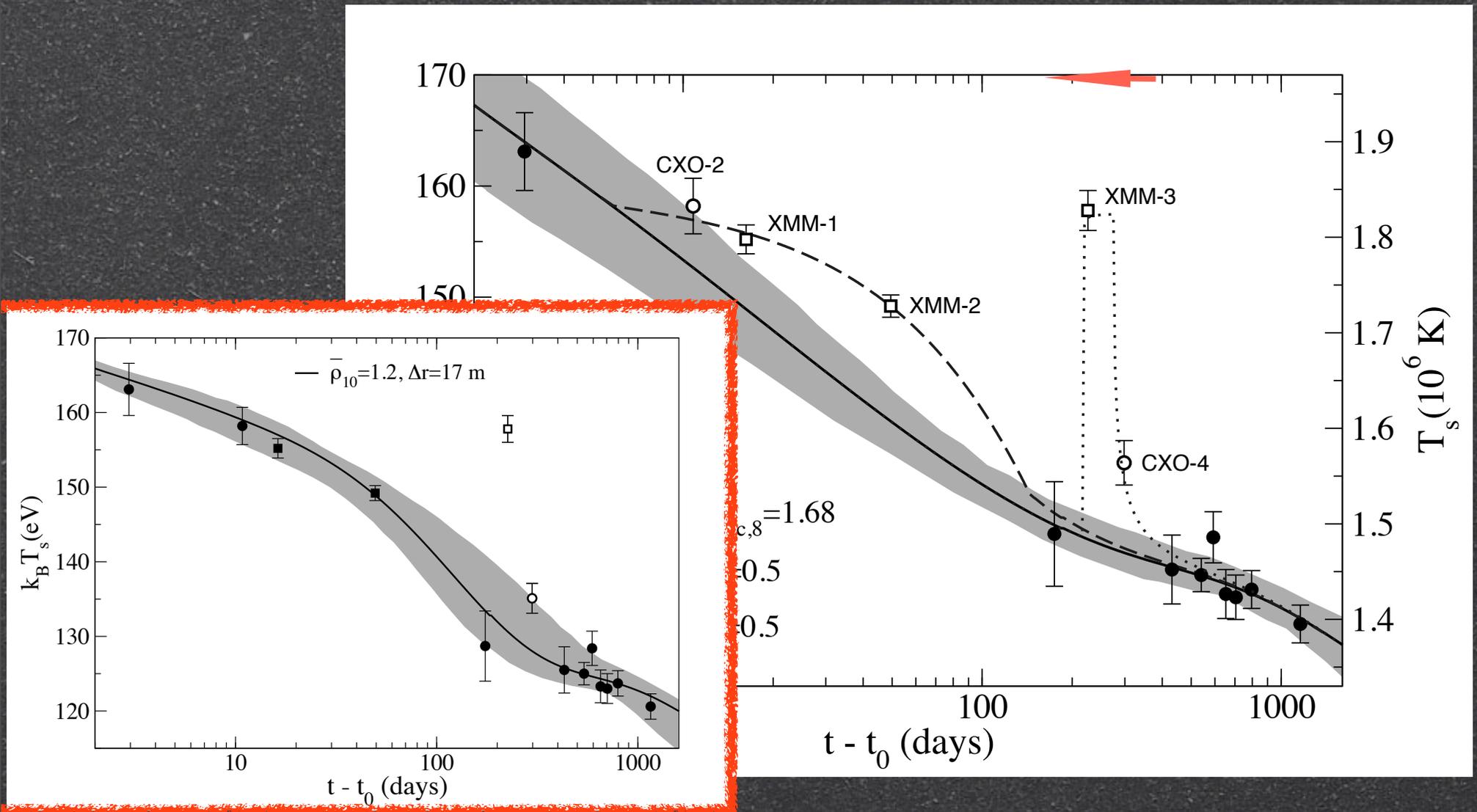


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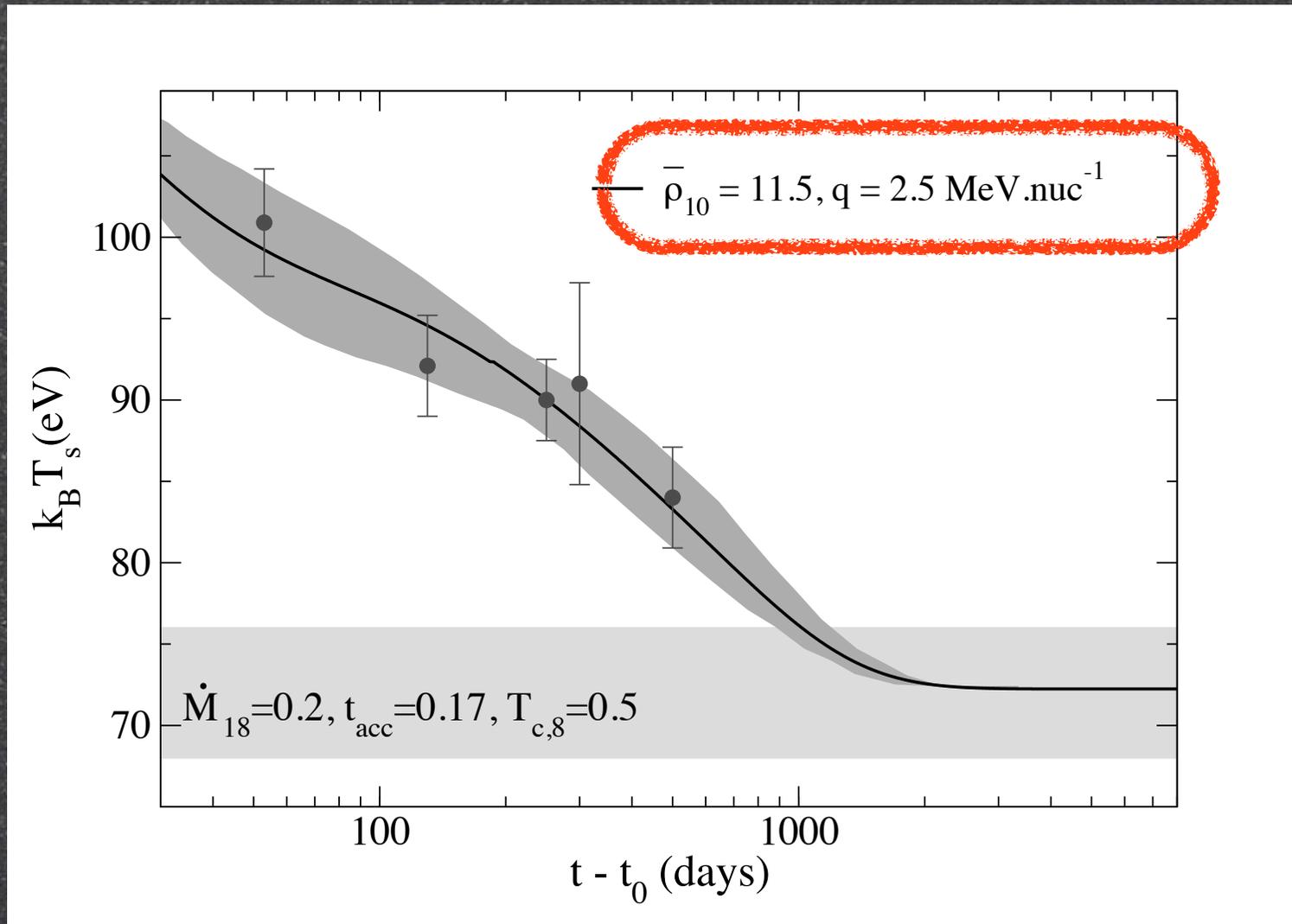
# XTE 1701-462: VARIABILITY

thermal isolation inner/outer crust?

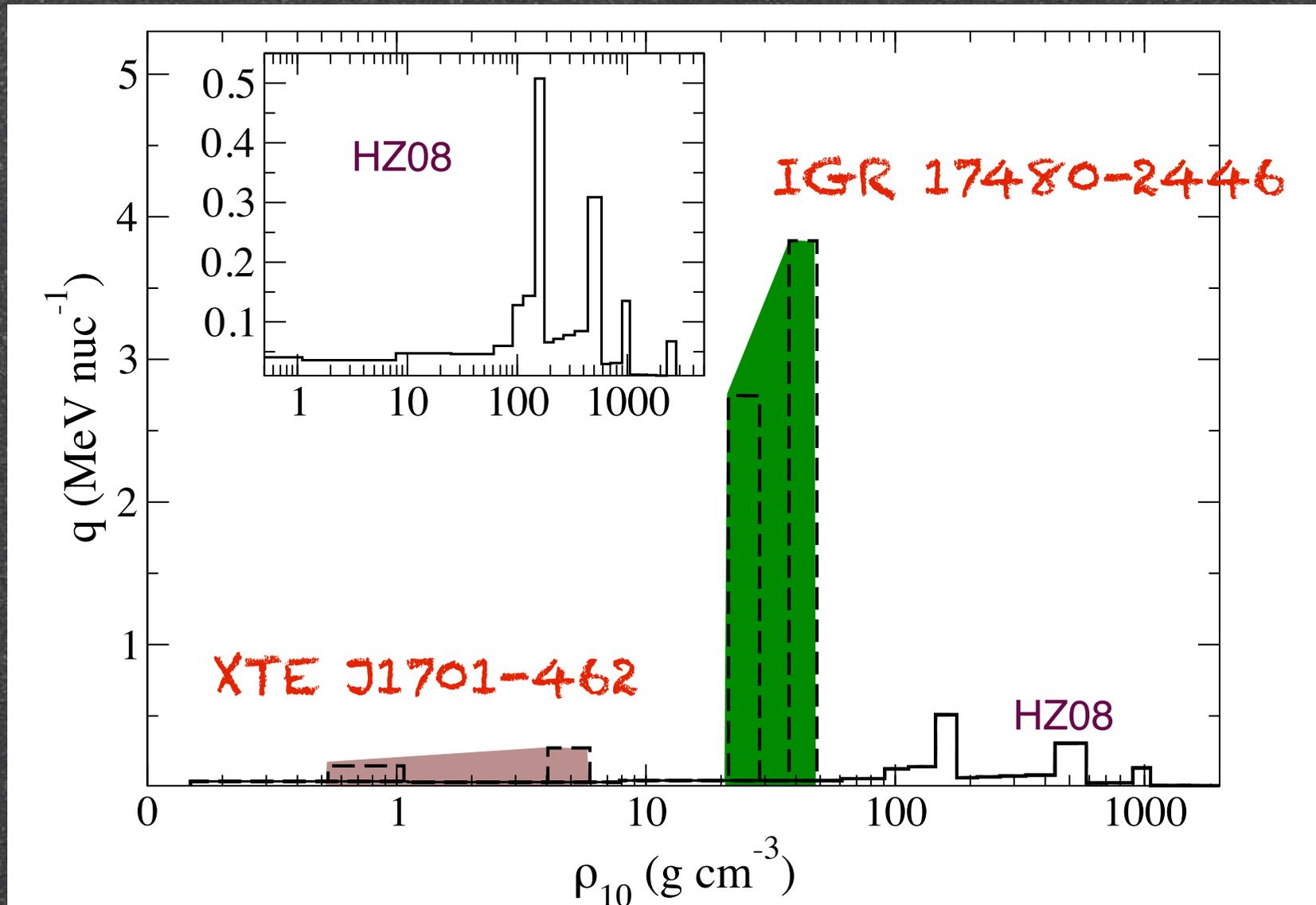


# IG J17480-2446: accreting only for 2 months

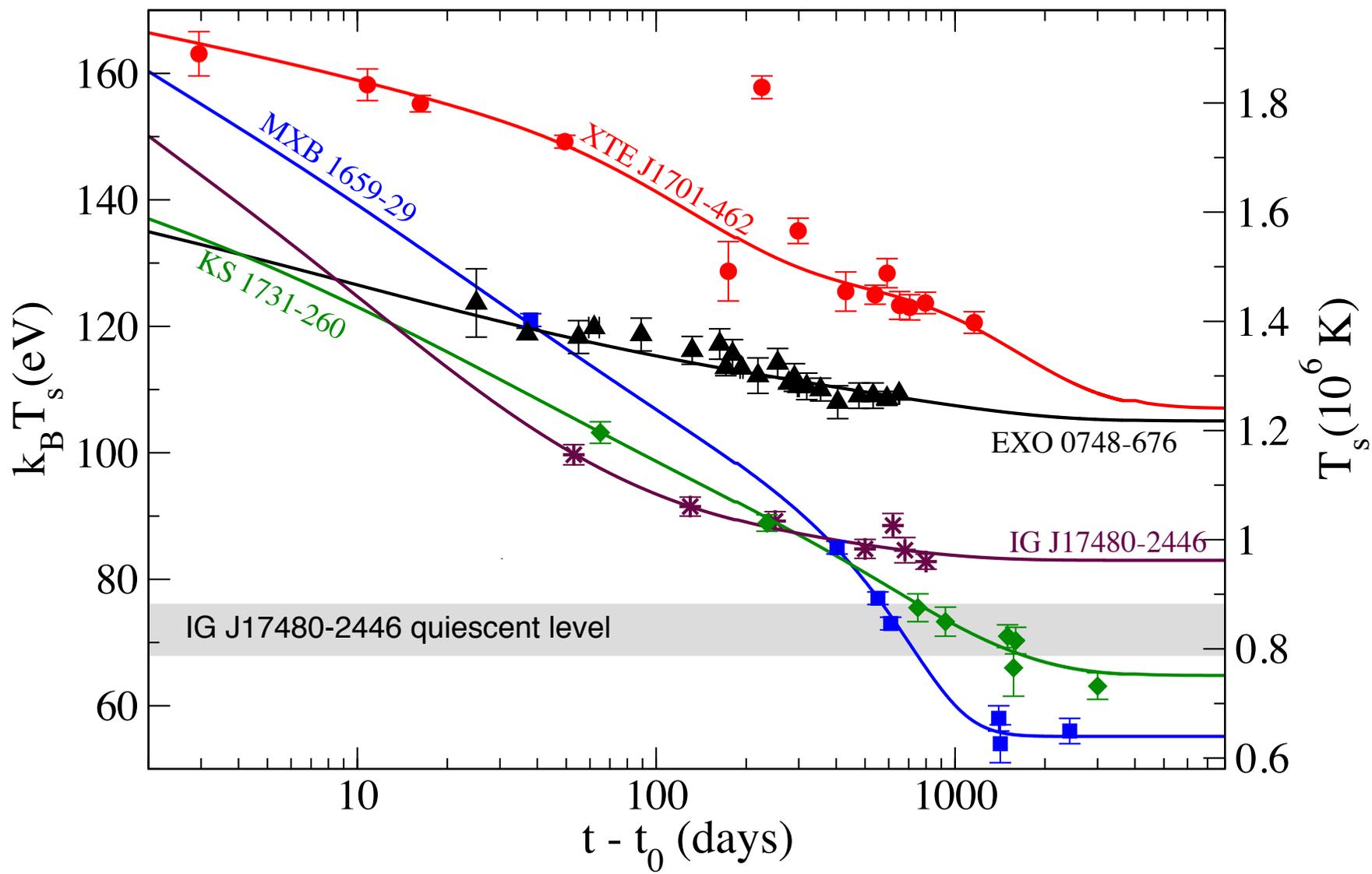
with additional heat sources



# Additional heat sources in the outer crust?



# Summary



# Conclusions

• Thermal evolution for the crust in quiescence for the five sources known

• Inclusion of the heated up envelope through the relation  $T_b(\dot{M})$

crust  $\leftrightarrow$  envelope  $\leftrightarrow$  outburst

• Constraints on

- crust microphysics ( $Q_{\text{imp}}$ , SF n gap)

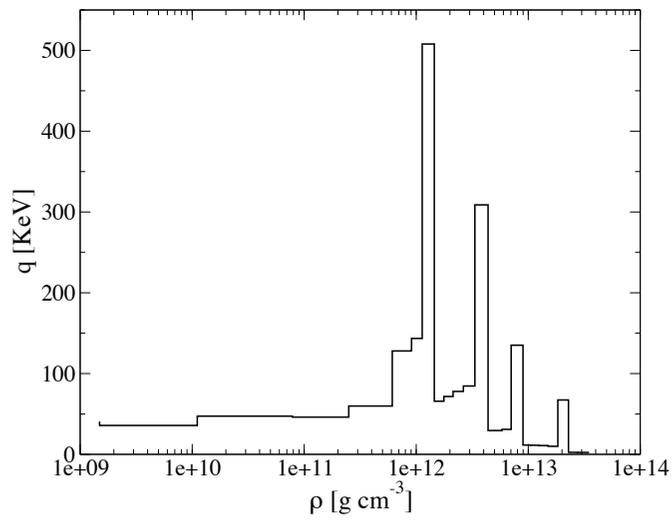
- envelope composition (H/He fraction)

• Three sources are "deep crustal coolers"

• Two sources are beyond crustal cooling

additional heat sources?, buried B, thermal wall?

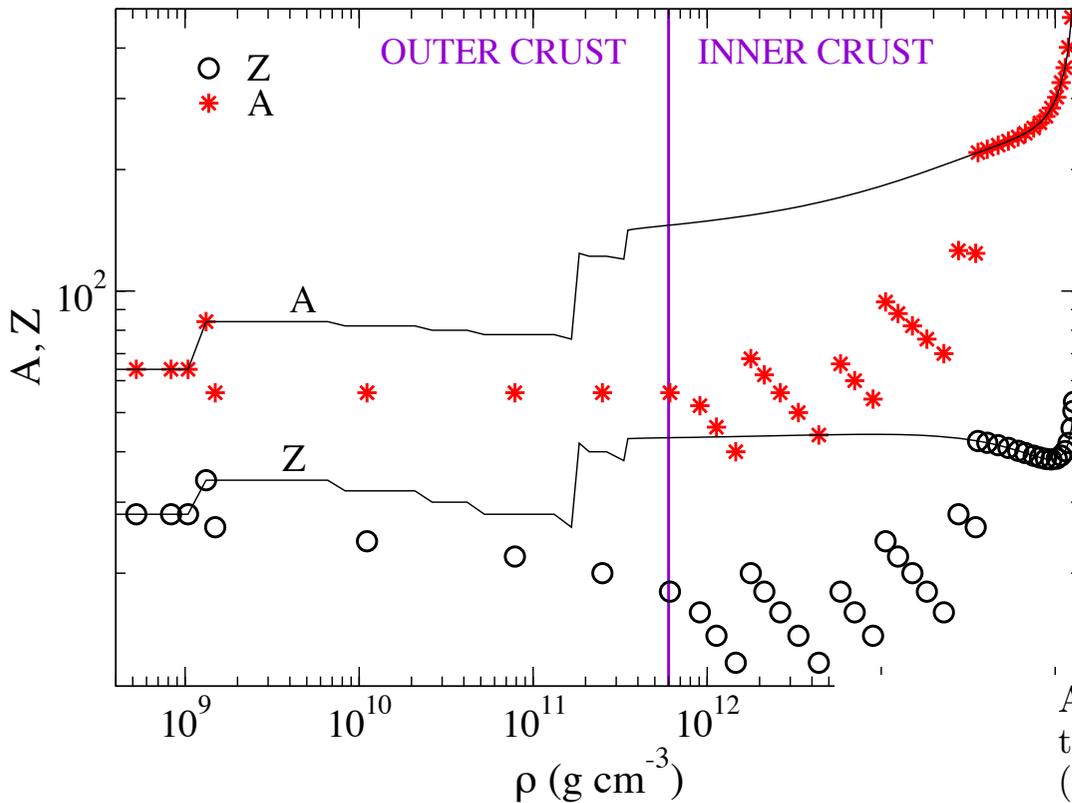
THANKS !



**Fig3. Left panel: Heat sources in the crust. Right panel: Non-equilibrium processes in the crust of an accreting neutron stars assuming that the X-ray ashes consist of pure  $^{56}\text{Fe}$ .  $P_j$  and  $\rho_j$  are pressure and density at which the reaction takes place. (Haensel & Zdunik 2008).**

$P$ (dyn cm $^{-2}$ )	$\rho$ (g cm $^{-3}$ )	Process
$7.23 \times 10^{26}$	$1.49 \times 10^9$	$^{56}\text{Fe} \rightarrow ^{56}\text{Cr} - 2e^- + 2\nu_e$
$9.57 \times 10^{27}$	$1.11 \times 10^{10}$	$^{56}\text{Cr} \rightarrow ^{56}\text{Ti} - 2e^- + 2\nu_e$
$1.15 \times 10^{29}$	$7.85 \times 10^{10}$	$^{56}\text{Ti} \rightarrow ^{56}\text{Ca} - 2e^- + 2\nu_e$
$4.75 \times 10^{29}$	$2.50 \times 10^{11}$	$^{56}\text{Ca} \rightarrow ^{56}\text{Ar} - 2e^- + 2\nu_e$
$1.36 \times 10^{30}$	$6.11 \times 10^{11}$	$^{56}\text{Ar} \rightarrow ^{52}\text{S} + 4n - 2e^- + 2\nu_e$
$1.980 \times 10^{30}$	$9.075 \times 10^{11}$	$^{52}\text{S} \rightarrow ^{46}\text{Si} + 6n - 2e^- + 2\nu_e$
$2.253 \times 10^{30}$	$1.131 \times 10^{12}$	$^{46}\text{Si} \rightarrow ^{40}\text{Mg} + 6n - 2e^- + 2\nu_e$
$2.637 \times 10^{30}$	$1.455 \times 10^{12}$	$^{40}\text{Mg} \rightarrow ^{34}\text{Ne} + 6n - 2e^- + 2\nu_e$ $^{34}\text{Ne} + ^{34}\text{Ne} \rightarrow ^{68}\text{Ca}$
$2.771 \times 10^{30}$	$1.766 \times 10^{12}$	$^{68}\text{Ca} \rightarrow ^{62}\text{Ar} + 6n - 2e^- + 2\nu_e$
$3.216 \times 10^{30}$	$2.134 \times 10^{12}$	$^{62}\text{Ar} \rightarrow ^{56}\text{S} + 6n - 2e^- + 2\nu_e$
$3.825 \times 10^{30}$	$2.634 \times 10^{12}$	$^{56}\text{S} \rightarrow ^{50}\text{Si} + 6n - 2e^- + 2\nu_e$
$4.699 \times 10^{30}$	$3.338 \times 10^{12}$	$^{50}\text{Si} \rightarrow ^{44}\text{Mg} + 6n - 2e^- + 2\nu_e$
$6.043 \times 10^{30}$	$4.379 \times 10^{12}$	$^{44}\text{Mg} \rightarrow ^{36}\text{Ne} + 8n - 2e^- + 2\nu_e$ $^{36}\text{Ne} + ^{36}\text{Ne} \rightarrow ^{72}\text{Ca}$ $^{72}\text{Ca} \rightarrow ^{66}\text{Ar} + 6n - 2e^- + 2\nu_e$
$7.233 \times 10^{30}$	$5.839 \times 10^{12}$	$^{66}\text{Ar} \rightarrow ^{60}\text{S} + 6n - 2e^- + 2\nu_e$
$9.238 \times 10^{30}$	$7.041 \times 10^{12}$	$^{60}\text{S} \rightarrow ^{54}\text{Si} + 6n - 2e^- + 2\nu_e$
$1.228 \times 10^{31}$	$8.980 \times 10^{12}$	$^{54}\text{Si} \rightarrow ^{48}\text{Mg} + 6n - 2e^- + 2\nu_e$ $^{48}\text{Mg} + ^{48}\text{Mg} \rightarrow ^{96}\text{Cr}$ $^{96}\text{Cr} \rightarrow ^{94}\text{Cr} + 2n$
$1.463 \times 10^{31}$	$1.057 \times 10^{13}$	$^{94}\text{Cr} \rightarrow ^{88}\text{Ti} + 6n - 2e^- + 2\nu_e$
$1.816 \times 10^{31}$	$1.254 \times 10^{13}$	$^{88}\text{Ti} \rightarrow ^{82}\text{Ca} + 6n - 2e^- + 2\nu_e$
$2.304 \times 10^{31}$	$1.506 \times 10^{13}$	$^{82}\text{Ca} \rightarrow ^{76}\text{Ar} + 6n - 2e^- + 2\nu_e$
$2.998 \times 10^{31}$	$1.838 \times 10^{13}$	$^{76}\text{Ar} \rightarrow ^{70}\text{S} + 6n - 2e^- + 2\nu_e$
$4.028 \times 10^{31}$	$2.287 \times 10^{13}$	$^{70}\text{S} \rightarrow ^{64}\text{Si} + 6n - 2e^- + 2\nu_e$ $^{64}\text{Si} + ^{64}\text{Si} \rightarrow ^{128}\text{Ni}$ $^{128}\text{Ni} \rightarrow ^{126}\text{Ni} + 2n$
$5.278 \times 10^{31}$	$2.784 \times 10^{13}$	$^{126}\text{Ni} \rightarrow ^{124}\text{Fe} + 2n - 2e^- + 2\nu_e$
$7.311 \times 10^{31}$	$3.493 \times 10^{13}$	$^{124}\text{Fe} \rightarrow ^{122}\text{Cr} + 2n - 2e^- + 2\nu_e$

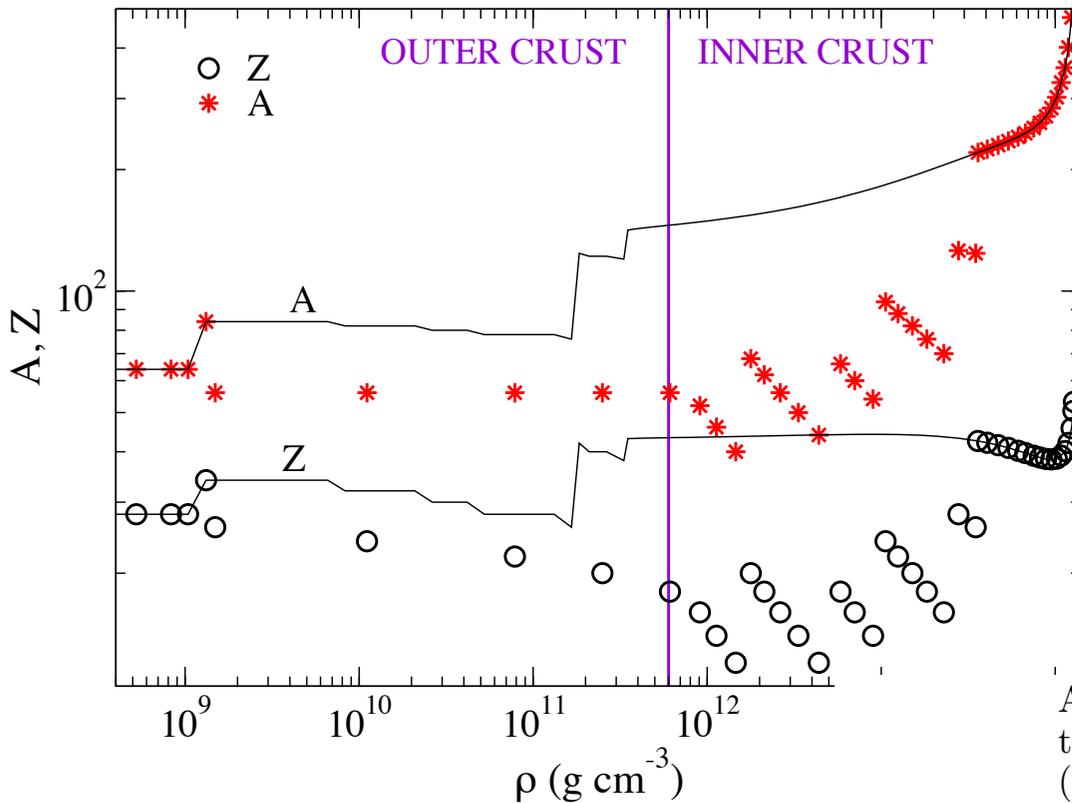
# Accreted crust composition



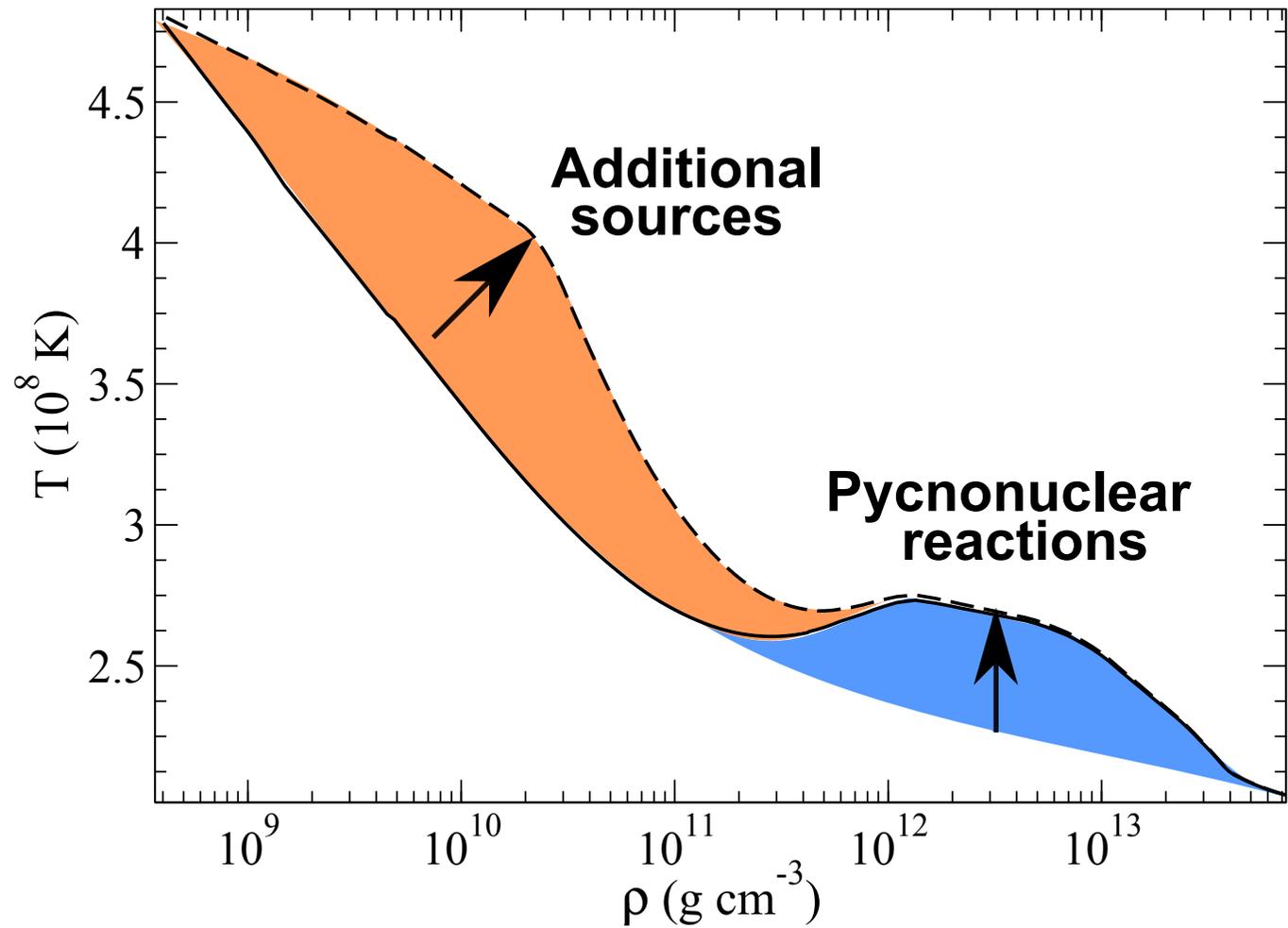
At low density we use the BBP Baym et al. (1971) equation of state. The crust-envelope interface is placed at  $(5 - 6) \times 10^8 \text{ g cm}^{-3}$  and we continue using BBP EoS to describe the crust up to the density  $1.49 \times 10^9 \text{ g cm}^{-3}$ . To take into account the effects of the accretion in the crust composition, we use the EoS presented in HZ08 in the range  $\rho = (1.49 \times 10^9 - 3.5 \times 10^{13}) \text{ g cm}^{-3}$ . This is a BBP-like EoS but modified by non-equilibrium nuclear reactions in the crust (see next Section). To describe the very high density region in the inner crust and the core we use Douchin & Haensel (2001), a Skyrme-type EoS which considers a nucleon-nucleon SLy effective interaction. For this chosen EoS the crust-core interface is at  $0.5 \rho_0$ , where  $\rho_0 = 2.8 \times 10^{14} \text{ g cm}^{-3}$  is the nuclear saturation density.

# Accreted crust composition

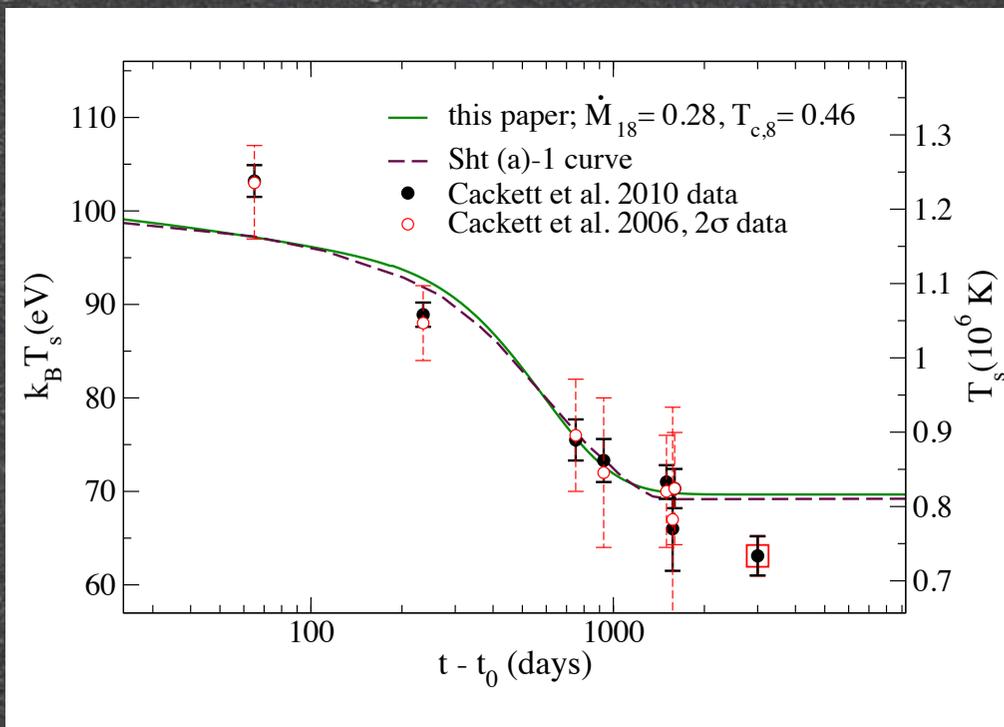
from Haensel & Zdunik 2008



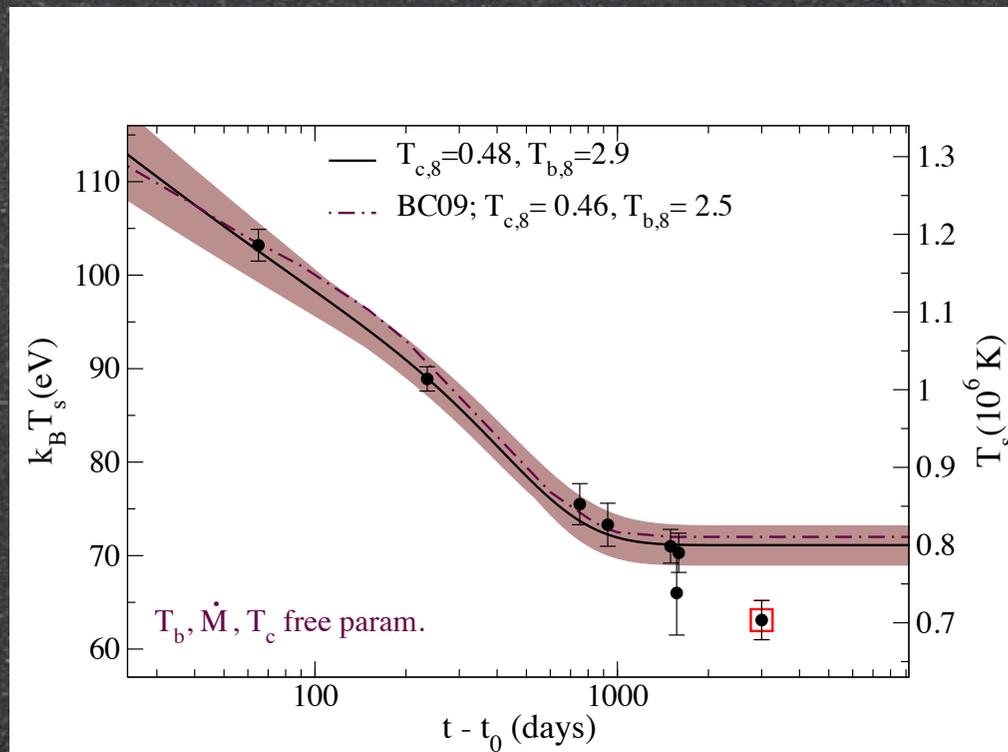
At low density we use the BBP Baym et al. (1971) equation of state. The crust-envelope interface is placed at  $(5 - 6) \times 10^8 \text{ g cm}^{-3}$  and we continue using BBP EoS to describe the crust up to the density  $1.49 \times 10^9 \text{ g cm}^{-3}$ . To take into account the effects of the accretion in the crust composition, we use the EoS presented in HZ08 in the range  $\rho = (1.49 \times 10^9 - 3.5 \times 10^{13}) \text{ g cm}^{-3}$ . This is a BBP-like EoS but modified by non-equilibrium nuclear reactions in the crust (see next Section). To describe the very high density region in the inner crust and the core we use Douchin & Haensel (2001), a Skyrme-type EoS which considers a nucleon-nucleon SLy effective interaction. For this chosen EoS the crust-core interface is at  $0.5 \rho_0$ , where  $\rho_0 = 2.8 \times 10^{14} \text{ g cm}^{-3}$  is the nuclear saturation density.

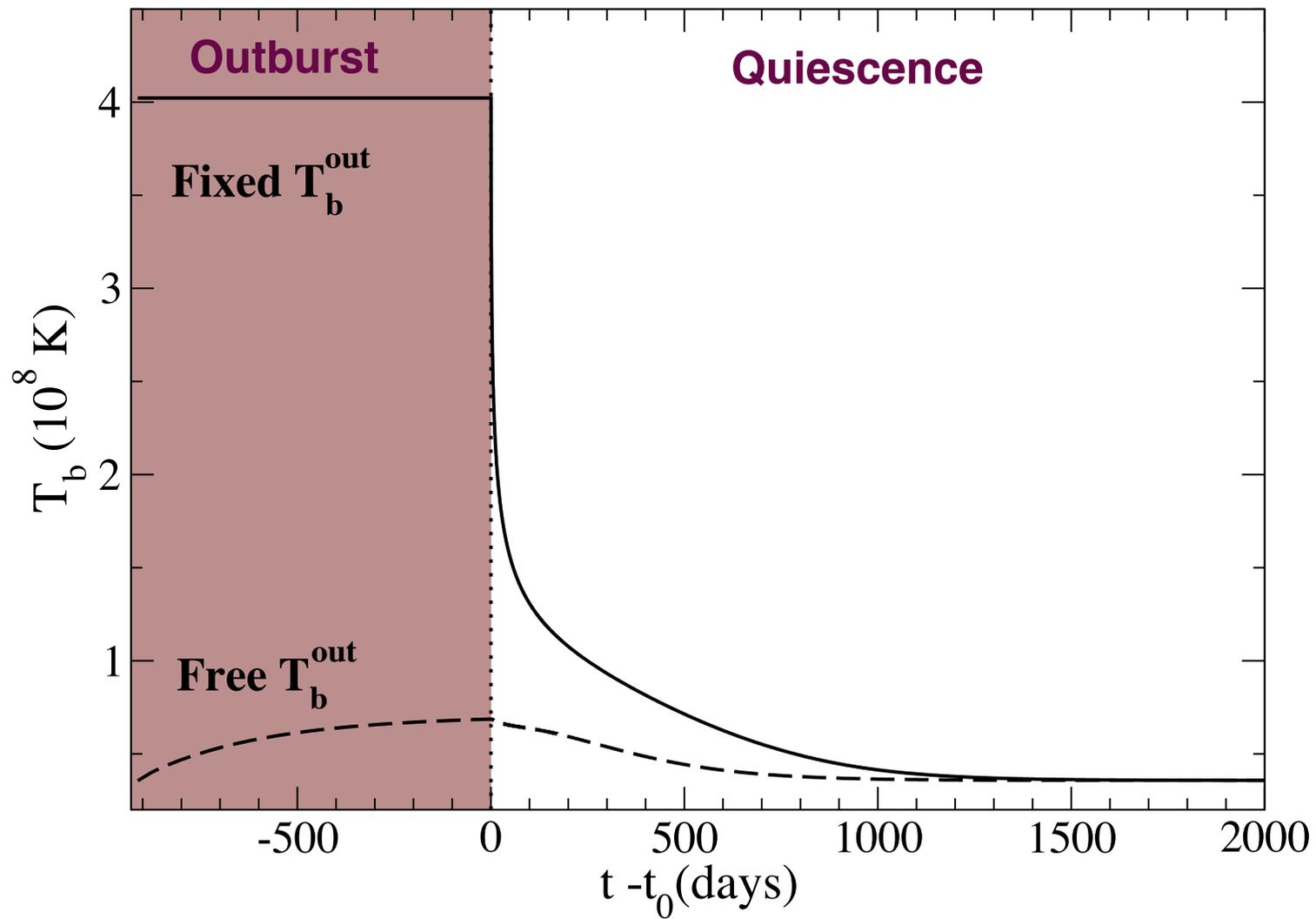


**Fig. 17.** Heat sources that influence on the initial thermal profile.



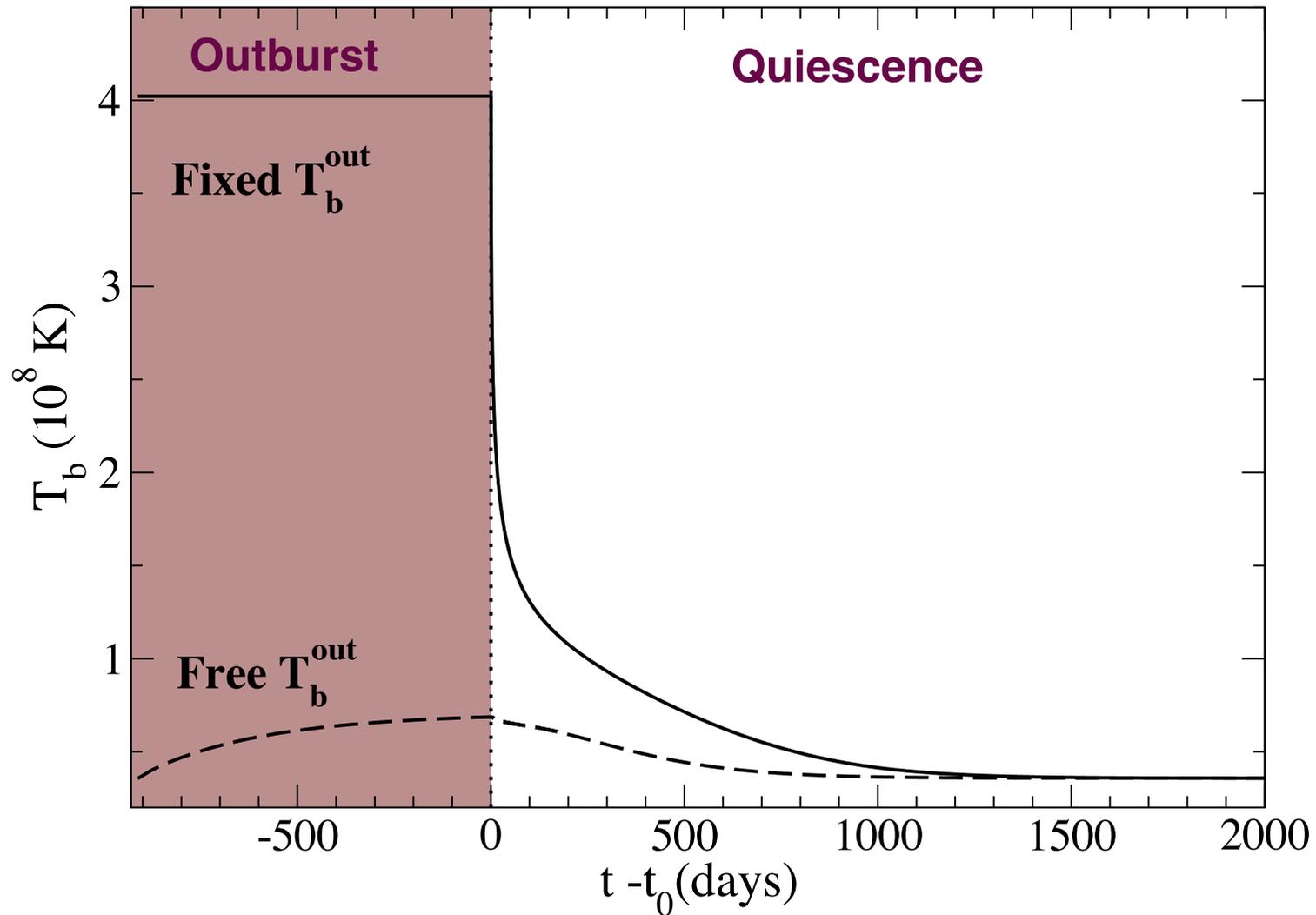
(a) Solid line is our curve with  $Q_{\text{imp}} = 2$ . Dashed curve is from Sht07.





(a) Temperatures at the crust-envelope interface

# The initial thermal profile



(a) Temperatures at the crust-envelope interface

