



On the energy of curved carbon surfaces

A multi-scale Density Functional Theory and Monte Carlo study

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19 September 2010

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▶ graphene, (single and multiwall) nanotubes, fullerenes

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Fullerenes treated with DFT

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graphene, (single and multiwall) nanotubes, fullerenes
 superior mechanical and electronic properties



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graphene, (single and multiwall) nanotubes, fullerenes

superior mechanical and electronic properties

▶ fullerene C₆₀ first dicovered in 1985 by Kroto *et al.*





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- Fullerenes treated Ring shell structure Monte Carlo

- graphene, (single and multiwall) nanotubes, fullerenes
- superior mechanical and electronic properties
- fullerene C_{60} first dicovered in 1985 by Kroto *et al.*
- carbon aggregates tend to be more stable in a form of onion-like structures



Banhart et al.. Phys. Chem. Lett. 269 (1997), 349.



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- Fullerenes treated with DFT Ring shell structure Monte Carlo simulations
- Description of the fullerenes
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- graphene, (single and multiwall) nanotubes, fullerenes
- superior mechanical and electronic properties
- ▶ fullerene C₆₀ first dicovered in 1985 by Kroto *et al.*
- carbon aggregates tend to be more stable in a form of onion-like structures
- possible use: supercapacitators, hydrogen storage



Banhart *et al.*, Phys. Chem. Lett. 269 (1997), 349.









• experiments suggest that the maximum number of layers is limited to $\approx 40 - 60$.

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 - ? Why?
 - ? May this be explained by mechanical instabilities occuring after a certain growth limit?

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- a continuum mechanics model of mechanical instability proposed by Bitsche *et al.* (submitted)







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 - input: surface stress pprox surface energy







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- ? Why?
- ? May this be explained by mechanical instabilities occuring after a certain growth limit?
- a continuum mechanics model of mechanical instability proposed by Bitsche *et al.* (submitted)
 - input: surface stress \approx surface energy
- the (excess) surface energy of a fullerene (or SWNT) is an increase of its total energy with respect to (planar) graphene





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Vienna Ab-initio Simulation Package with GGA-PAW and LDA-US pseudopotentials



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 only pentagons and hexagons ~ exactly 12 pentagons, variable number of hexagons (Euler's formula)

(# of corners) – (# of edges) + (# of facets) = 2





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▶ focus on fullerenes with icosahedral symmetry (with the exception of C₇₀)



Fullerenes from DFT



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parts of fullerene surfaces with dangling bonds "saturated" with H atoms

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(a) one-ring shell (b) two-ring shell (c) three-ring shell Ring shell structure

- parts of fullerene surfaces with dangling bonds "saturated" with H atoms
- C–H bond length optimised to fit best the graphene energy (per bond)

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- parts of fullerene surfaces with dangling bonds "saturated" with H atoms
- C-H bond length optimised to fit best the graphene energy (per bond)
- +/- correspond to spherical structures
 - contain no pentagons
 - $+\,$ can easily access any desired radius

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construction of classical inter-atomic potentials "coherent"
with the present DFT calculations
 stretching: Morse potential

$$E^{S}(r_{ij}) = E_{0}\left[\left(1 - \exp(\beta(r_{ij} - r_{0}))\right)^{2} - 1\right]$$



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construction of classical inter-atomic potentials "coherent" with the present DFT calculations

- stretching: Morse potential
- bending: harmonic potential

$$E^{B}(heta_{ijk}) = rac{1}{2}k_{ heta}\left(\cos heta_{ijk} - \cos heta_{0}
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torsion

$$E^{T}(\phi_{ijkl}) = rac{1}{2}k_{\phi}\left(1 - \cos 2\phi_{ijkl}
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 DFT and MC give the same exponent but different off-sets (power-law fit: E = E₀ × R^β)







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- ring shell models fail to describe individual fullerenes but may be appropriate for (spherical) onion-like structures



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 geometry and energy of carbon fullerenes were studied by a combination of the density functional theory and Monte Carlo methods

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- geometry and energy of carbon fullerenes were studied by a combination of the density functional theory and Monte Carlo methods
- DFT-compatible inter-atomic potentials for MC were constructed

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- curvature induced excess surface energy follows a power law as a function of the structure mean radius

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- curvature induced excess surface energy follows a power law as a function of the structure mean radius
- power law exponent for the surface energy per bond: $\beta^{\mathsf{DFT}} = -1.43$ and $\beta^{\mathsf{MC}} = -1.40$

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- \blacktriangleright perfect spherical structures using approximative shell models yielded $\beta\approx-2.5$
- only stretching (two-particle) and bending (three-particle) interactons are not able to reproduce the correct fullerene geometries (i.e. bond-length and angle distributions leading to extended flattened areas), for which at least the four-particle interaction (i.e. torsion) has to be included.

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Thank you for your attention!

Financial support by the START Program (Y371) of the Austrian Science Fund (FWF) is greatly acknowledged.



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